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- **1** Spatially-microstructured topology of
- 2 chalcogenide glasses by a combination of
- 3 electrothermal process and selective etching
- 4 for functional infrared media
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14 Abstract: Chalcogenide glasses exhibit a wide transparency domain spanning from near 15 infrared (IR) to mid-IR and thus, have become highly attractive optical materials in a range of 16 applications. Controlling the topology of these glasses can be seen as a key aspect for the design of optical elements such as gratings, metasurfaces, waveguides, and other diverse refractive 17 18 and diffractive optical components. Here, we demonstrate the structuring of large, millimeter 19 square areas that have been structured at the micrometer scale employing an easy two-step 20 process, incorporating a micro-poling step followed by immersion in an amine solvent. Ge-Sb-21 S-Na glasses have been investigated, and the influence of the sulphur and sodium content on 22 the pre- and post-poling material dissolution response has been discussed. Three compositions 23 of varying sulphur and sodium content were selected to study the influence of thermal poling 24 using either a homogeneous or a structured electrode. It was found that either a large difference 25 in dissolution rates of poled and unmodified regions or a large poled layer thickness leads to 26 the generation of significant topological contrast. The origin of the poled region's selective 27 etching has been explained on the basis of a poling-induced density decrease. Finally, it was 28 demonstrated that when the targeted resolution is micrometric, this rather easy process could 29 be employed as an alternative to classical lithography techniques.

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32 1. Introduction

33 The control of a material's surface topology is a key aspect of its design and potential use in 34 optical applications. In particular, a wavelength length-scale (nano- or micrometer length-scale) 35 control is necessary to produce diffractive optical elements, [1] antireflective surfaces, [2] 36 waveguides, [3] plasmonic surfaces [4] or metasurfaces. [5,6] Multiple processes have been 37 shown to be suitable to structure surfaces, including lithographic techniques, which are limited 38 to the wavelength of the tool. These methods, inherited from the microelectronic domain are 39 among the most common and have been recently applied to low loss hybrid 40 (diffractive/refractive) optical elements and metastructures based on chalcogenide phase 41 change alloys. [7–9] Lithography is a cumbersome method employing multi-step processes 42 including deposition and patterning of a photosensitive sacrificial layer, which subsequently 43 allows the selective destruction of the material of interest's surface to realize the target 44 structure. In each case, material-specific protocols are required, while the determination of such

45 a 'recipe' is time consuming, and requires compatibility between resist, etchant, and writing46 wavelengths for each desired composition and element's design.

47 In the 1990's, nanoimprint lithographic techniques were developed, [10] allowing generation 48 of low-cost, sub-25 nm structures over large areas $(15 \times 18 \text{ mm}^2)$. These techniques differ from 49 classical lithography since a mold is employed to directly structure the desired material 50 (generally polymeric resins). The mold, pressed against the sample (under load at temperature), 51 generates a thickness contrast, and the compressed zones are preferentially etched or displaced 52 for subsequent etching using either a wet or dry reactive ion etch step. Employing this 53 technique, Kohoutek et al. managed to periodically structure the surface of an arsenic selenium 54 chalcogenide glass resulting in an effective diffraction grating. [11] Here, the authors had to 55 heat up their system to 5 °C below the glass' transition temperature (T_o) (i.e., 225 °C) in order to imprint the pattern. However, for glasses exhibiting higher T_g , the required pressure and 56 57 temperature to imprint the patterns increase and can cause several problems, including higher 58 processing costs, difficulty to realize the process over large areas, material volatilization, and 59 reduced life expectancy of the molds. [12–14] To circumvent this difficulty. Takagi et al. 60 applied an electric field and used their mold as an anode. [12] While the applied voltage in this 61 case was only of about a hundred of volts, and the pressure was one of the driving force, the 62 nanoimprint process employed by the authors is similar to thermal micro-poling. The thermal 63 poling method consists of heating an amorphous material below its T_g while applying a strong 64 electric field (several kilovolts per millimeter) with a structured electrode. [15] The sample is 65 then brought back to room temperature before removal of the electric field. Brunkov et al. and 66 Lipovskii et al. demonstrated a control of silicate and nano-composite glasses' surface 67 topologies using this process. [14,16,17] However, the reliefs obtained are only (typically) tens 68 of nanometers high which is insufficient to consider the fabrication of efficient gratings for 69 instance. Ikutame et al. suggested augmenting the poling treatment with one inducing chemical 70 attack. Here, they obtained relief heights higher with subsequent chemical attack than with 71 poling alone. [18]

72 The first studies regarding chemical attacks of poled glasses date back to the early days of 73 thermal poling. In 1974, Carlson noticed that the poled surface (hence depleted from its alkali 74 cations) of a silicate glass was shown to be more resistant to molten alkali salts attack than the 75 untreated surface. [19] Others relied on the difference in etching rate difference of the poled 76 and unpoled zones to measure the poled thickness in fused silica and silicate glasses. [20,21] It 77 is only more recently that authors started focusing on post-poling chemical attack in the optic 78 of realizing devices such as diffraction gratings employing structured electrodes. [22] Poled 79 silicate glasses are selectively dissolved by acidic attack, [23,24] but authors also demonstrated 80 a dry etching route by plasma or reactive ion etching where preferential removal rates differ 81 between poled and unpoled regions. [24,25] These two possibilities allow one to preferentially 82 etch either the unpoled or the poled zones. It follows that the electrode's pattern can be either 83 directly transferred on the glass' surface or negatively transferred. The step's heights obtained 84 by these techniques on the (to-date) studied oxide glasses' family (silicate glasses) are measured 85 to be between 0.5 and 1.2 μ m.

86 Dissolution of chalcogenide glasses (ChGs) has been the subject of many studies since the 87 1980's. Numerous works focus on thin films synthesis by dissolution of bulk glasses followed 88 by spin coating. [26–29] It was demonstrated in the last decade that sulfur glasses were 89 effectively dissolved in amine solvents such as ethylene diamine (EDA). Furthermore, our 90 previous works on G-Sb-S-Na glasses have shown the efficiency of micro-poling on these 91 glassy compositions. Indeed, after the treatment, the sodium repartition, the structural 92 modifications, the index variations, and the second harmonic optical response follow the pattern 93 of the structured electrode. [30–32] In the present work, our aim has been to transfer the know-94 how realized in post-poled selective dissolution of silicate [22-24] to this chalcogenide glass 95 system while identifying and quantifying the key parameters that can control the topologies of

96 the resulting surfaces.

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98 2. Experimental method

99 2.1 Glass synthesis

100 Glasses were prepared using high purity elemental Ge, Sb, and S (Alfa Aesar 99.999%). 101 Anhydrous sodium sulfide (Na₂S – purity unspecified) was used to incorporate sodium in the 102 glass matrix. The nominal composition of the glasses can be written as: $(100-x)(Ge_{25-y}Sb_{10}S_{65+y})$ 103 + xNa with x = 0, 2, 5.5 and y = 0, 2.5, 5, 7.5. Two series of compositions were selected. (1) 104 The first one with varying S to Ge ratio (i.e., increasing amount of sulphur) and fixed sodium 105 content: A-Na₂, B-Na₂, C-Na₂, D-Na₂ where x = 2 and y = 0, 2.5, 5, 7.5, respectively. (2) The 106 second one, with varying sodium content and fixed S to Ge ratio: C-Na₀, C-Na₂, C-Na₅₅ where 107 x = 0, 2, 5.5, and y = 5, respectively. These different parameters are summarized in the Table 1. 108 Table 1. The S/Ge ratio, the sodium content, the densities, and the glass transition temperatures of the



synthetized	glasses as well:	as the name they	are referred to.

Sample	S/Ge ratio	Na content (at.%)	Density (g/cm ³)	T_g (°C)
A-Na ₂	2.6	2	3.163 ± 0.006	313 ± 2
B-Na ₂	3.0	2	3.099 ± 0.006	270 ± 2
C-Na ₂	3.5	2	3.057 ± 0.012	241 ± 2
D-Na ₂	4.1	2	2.982 ± 0.006	-
C-Na ₀	3.5	0	3.06 ± 0.01	274 ± 2
C-Na _{5.5}	3.5	5.5	3.05 ± 0.01	250 ± 5

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Raw materials for 15 g batches were weighed out in a glove box under nitrogen atmosphere and inserted in a quartz ampule. The system was put under vacuum (10^{-2} mbar) and sealed using an oxygen-methane torch. The sample was placed in a rocking furnace heated at a rate of 1 °C/min up to 850 °C for 12 h. After this time period, rocking was stopped, and the temperature decreased to 750 °C before quenching in water. The sample was then placed into an annealing furnace for 6 h at 40 °C below its glass transition temperature (T_g). Finally, glasses were cut and polished to obtain 1 mm thick disks of 1 cm diameter. The glass transition temperatures were determined by differential scanning calorimetry (Netzsch DSC 204 F1 Phoenix) using aluminum pans and a heating rate of 10 °C/min up to 550 °C. Glass densities were determined using the Archimedes method by immersing the samples in diethyl phthalate at room temperature. The measured T_g and densities are reported in the Table 1.

122 2.2 Vibrational spectroscopy

123Raman spectroscopy to assess structural changes was carried out and two-dimensional spatial124Raman intensity maps were acquired using a LabRAM HR Evolution (Horiba) spectrometer125with a laser source at 785 nm and a $100 \times (N.A.= 0.9)$ objective with a one micron resolution.

126 2.3 Thermal poling

127 Thermal poling was performed under flowing nitrogen. The 1 cm diameter glass slides were 128 heated up to 170 °C at a 15 °C/min rate (temperature fixed for all samples). A DC voltage of 129 3 kV was then applied (at a 320 V/min rate) by a 5×5 mm² electrode placed at the center of the glass disk and left on for 30 min. The samples were then brought back to room temperature 130 131 before removal of the DC field. Two types of anodes were employed in this work. They 132 consisted of a 100 nm platinum film deposited on a glass slide ($5 \times 5 \times 1 \text{ mm}^3$) and were either 133 (1) left as such (homogeneous electrodes that are everywhere conductive) or (2) 134 lithographically structured to form $(40 \times 40 \ \mu\text{m}^2)$ non-conductive squared patterns delimited 135 by 10 µm large conductive platinum grid. On the cathode side, a piece of silicon wafer was

used, and a microscope coverslip was placed between the silicon and the sample to preservethe optical quality of the sample on this side.

138 2.4 Dissolution

139 Throughout this work, the solvent where the dissolution was carried out was EDA mixed at 140 15 volume % in dimethyl sulfoxide (DMSO). The dissolution of samples (i) prior to poling, (ii) 141 after a homogenous poling, and (iii) after a structured poling were studied. For these three 142 situations, different methods had to be developed to follow the dissolution rate. Specifically, 143 (i) In order to track the dissolution rate of virgin glasses, we chose to look at the dissolved mass 144 of the different samples with respect to the immersed time in the solvent. The one-centimeter 145 square and 1 mm thick glass samples were weighted out with a 0.01 mg precision scale. They 146 were immersed in the solvent and placed in an ultrasound bath for a given time. They were then 147 rinsed with distilled water and dried before being weighted again and placed back to the 148 ultrasound bath. (ii) To track the dissolution of the poled layer, weighting the samples is not 149 adapted anymore as this layer is only several micrometer thick. Drawing on the methodology adopted by Reduto et al., [23] we chose to protect a part of the surface with a metallic (gold) 150 151 thin film (Q300TD, Quorum Technologies) and to measure the step height between the metallic 152 film (not affected by the solvent) and the poled glass' surface with respect to the immersed time 153 in the solvent. These measurements were also performed on virgin (non-poled) gold-protected 154 samples for reference. The height step was measured with a Bruker's Dektak 6M stylus profiler. 155 The measured height corresponded to the mean value of at least four measurements at each 156 immersion time. Unlike the dissolution of the virgin glasses, the ultrasound bath was not used here to prevent the deterioration of the gold thin film. (iii) The samples poled with the structured 157 158 electrode were immersed in the solvent in an ultrasound bath for different times. They were 159 rinsed with distilled water, and dried. The topological profiles around the imprinted structures 160 were then measured by atomic force microscopy (AFM) on an AFM Dimensions Icon Bruker 161 instrument in PF-QNM mode. The topographic images were acquired using a ZYGO 6300 162 white-light interferometer.

163 3. Results

164 3.1 Dissolution of the virgin glasses

165 We first discuss the effect of the glasses' compositions on the dissolution rate of the samples 166 prior to poling. In Fig. 1 are presented the dissolved masses with respect to the immersed time 167 in the solvent. Fig. 1a illustrates the effect of varying sulphur content (for a fixed sodium content) while in Fig. 1b, the effect of the sodium content is investigated. On the first hand, 168 169 from Fig. 1a, one can clearly observe that the higher the sulphur content is, the higher the 170 dissolution rate is. The glass D-Na₂ is dissolved about four times faster than the stoichiometric 171 glass (A-Na₂). On the other hand, the effect of the sodium content (Fig. 1b) is less pronounced. 172 Taking into account the error bars, we can consider that, in the sodium concentration range 173 evaluated in this study, the effect on the dissolution rate is negligible.



Fig. 1. Virgin glasses' dissolved mass with respect to time as a function of (a) the S/Ge ratio and of (b) the sodium content for 1 cm² area samples. The error was estimated to 25 mg to account for the experimental error coming from the weighting and the sample's area estimation, the lines correspond to linear regressions.

179 3.2 Dissolution of homogenously poled glasses

To highlight the effect of poling and the resulting imparted change in the glass on the dissolution rate, the adopted methodology, as described in the experimental section, was to protect a part of the sample with a gold film and to measure the step height between this film and the glass as a function of the immersed time. Three compositions possessing two different S/Ge ratio and two different sodium contents – namely A-Na₂, C-Na₂ and C-Na_{5.5} – were selected for evaluation. The results of this set of experiments are presented in Fig. 2. One can first notice a linear increase in the step height for the virgin glasses that enables the dissolution rate (v_{UP}) of virgin-like (unpoled) glass regions to be estimated. These values are gathered in the Table 2. From this table, it appears that C-Na_{5.5} has the fastest dissolution rate, followed by C-Na₂ and A-Na₂. These observations are consistent with the dissolved mass evolutions seen in Fig. 1. In the poled samples, one can clearly identify two dissolution rate reaches that of the virgin glass. Here, the slope inflection point corresponds to the thickness of the poled layer, L_{poled} . For each sample, its value as well as the dissolution rate of the poled region (v_P) are reported in the Table 2.

195 In order to generate structures with a maximum contrast, it is crucial to consider Δh_{max} , the 196 maximum gap between the poled and the unpoled zones. This parameter is illustrated by the 197 schematic in Fig. 2d. It is formulated as Eq. (1):

$$\Delta h_{max} = L_{poled} - L_1 = L_{poled} - t_1 v_{UP} = L_{poled} - L_{poled} \frac{v_{UP}}{v_P}$$
$$= L_{poled} \left(1 - \frac{v_{UP}}{v_P}\right) \#(1)$$

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From Eq.(1) and Fig. 2d, it follows that there are two methods to maximize Δh_{max} : having a sample exhibiting either a thick poled layer, or a large difference in dissolution rates of the poled and unpoled regions. In the selected compositions, two good candidates stand out: C-Na₂ with its eight microns-thick poled layer and C-Na_{5.5} with a significant difference in dissolution rates of the poled and unpoled regions (lowest ratio $v_{\text{UP}}/v_{\text{P}}$, see Table 2).



Fig. 2. Step heights measured between the protective gold film and the glass' surface as a function of the immersion time in the solvent for the compositions (a) A-Na₂, (b) C-Na₂, and (c) C-Na_{5.5}. The dots and lines correspond to the measured points and linear regressions, respectively. The error is mainely linked to the use of the profilometerand was estimated to be between 300 and 500 nm. The red lines and points correspond to the virgin glasses while the black ones correspond to the homogeneously poled samples. A schematic (d) illustrates the key parameters namely: L_{poled} , the poled layer thickness, v_P and v_{UP} the poled and unpoled dissolution speeds, and Δh_{max} , the maximum height that can be reached between the poled and the unpoled zones.

 Table 2. Key parameters extracted from Fig. 2: poled zone's thickness, poled and unpoled regions' dissolution speeds, and their ratio as well as the theoretical maximal height that can be reached for each sample calculated from Eq. (1).

Sample	$L_{\text{poled}}(\mu m)$	$v_{\rm P} ({\rm nm/s})$	$v_{\rm UP} ({\rm nm/s})$	$v_{\rm UP}/v_{\rm P}$	$\Delta h_{\max} \left(\mu m \right)$
A-Na ₂	2.4 ± 0.3	4.7 ± 0.5	2.2 ± 0.2	0.47	1.3 ± 0.2
C-Na ₂	8.0 ± 0.8	8.5 ± 0.9	4.7 ± 0.5	0.55	3.6 ± 0.4
C-Na _{5.5}	4.3 ± 0.4	14.7 ± 1.5	5.8 ± 0.6	0.39	2.6 ± 0.3

217 3.3 Dissolution of the micro-poled glasses

In this section, we no longer use a homogenous, two dimensional conductive electrode but an electrode that has been structured such that it allows for the micro-scale control of the poling process and opens the way for the design of surfaces for specific applications such as gratings. [22,24] As shown in prior work on this particular glassy system, the size and shape of the electrode influences the spatial extent of the electric field that serves to cause migration of the charged alkali ion (here sodium). [30,32] An optical image of the structured electrode used in this study is presented in Fig. 3a. It consists of a ten micron wide metallic grid separated by $40 \times 40 \ \mu\text{m}^2$ squares that are not conductive. Here, we define two regions on the micro-poled sample: the region in contact (IC) with the metallic grid during the process and the region not in contact (NIC) with the grid.

228 Following the methodology adopted in our previous works, [31,33] the effect of micro-poling 229 (prior to any dissolution) on the glass' structure can be readily investigated by Raman microscopy. These results are presented for the C-Na_{5.5} sample throughout Fig. 3. The Raman 230 231 spectra extracted from the NIC and IC with the grid regions and are presented in the top part of 232 Fig. 3b. Since only slight differences between these two spectra can be observed, the Raman 233 difference spectrum (IC minus NIC) is also presented in Fig. 3b. The spatial evolution of the 234 Raman response is illustrated by the cartographies of the five bands selected on the Raman 235 difference spectrum.

236 In our previous work, a complete attribution of these bands is presented relying mainly on the 237 works of Ward, [34] Lucovsky et al. [35] and Koudelka et al. [36] Specifically, in these glasses, 238 antimony and sulphur form $SbS_{3/2}$ pyramids, while germanium is organized as $GeS_{4/2}$ 239 tetrahedra. In glasses without sodium, the Sb-S stretching modes appear at 302 cm⁻¹ [band (ii)]. 240 When sodium is added, they shift to lower wavenumbers [band (i)]. The GeS_{4/2} tetrahedra are mainly organized as corner sharing (CS) configuration [band (iii)] but also in edge sharing (ES) 241 242 configuration [band (iv)]. Sodium-rich glasses present more ES units than sodium poor glasses, 243 the opposite goes for CS units. As for the band (v), it is attributed to S-S homopolar bonds in 244 S_n chains or S_8 rings.

In the present work, similar patterns are observed for all Raman intensity maps. In particular, sodium-rich glassy structures exhibit bands such as (i) and (iv) which are negative in the IC zones and positive in the NIC regions. The opposite behavior is seen for the sodium-poor glassy structures related bands such as bands (ii) and (iii). This conclusive variation in network configuration confirms the effectiveness of the micro-poling process that results in the micro-scale control of the structural variations as well as the observed sodium content distribution within that network.

252 The spatial intensity profile of the band (i) is shown in Fig. 3d. On this graph, the region 253 between the dotted lines corresponds to the region IC defined by the metallic grid. As previously stated, this band is negative in this region but one can also notice a gradient in the 254 255 intensity profile. The structural modifications are thus not only effective where the metallic 256 electrode is in contact with the glass surface but they are also extend to adjacent regions to 257 about five microns on each sides of the grid. In other words, there is not a perfect overlap between the IC and poled zones as well as between the NIC and unpoled regions: the poled 258 259 zone extends beyond the IC zone. This is the result of an in-plane poling similar to that observed 260 in other systems, [15,31,37,38] We have shown that it is possible to take advantage of this 261 gradient in our previous work where we specifically aimed to design flat micro lenses through 262 inducing a gradient in refractive index within this glassy system. [30]

263 Now that the microscale patterning of the structural variations generated by the micro-poling is confirmed, we now focus on the dissolution behavior of this sample. The topology profiles 264 265 measured by AFM at different dissolution times are presented on the Fig. 3e. Immediately after 266 poling (at t = 0 s), the sample's surface is flat. After immersion in the solvent, the IC zone can 267 be seen to be selectively etched. This is consistent with the results from Fig. 2 where the poled 268 samples exhibited more rapid dissolution than the virgin glasses. After 60 s of immersion, the 269 glass surface forms a 1.7 µm high step; after 300 s, it reaches 3.7 µm. Comparing the AFM 270 profiles with the Raman intensity profile from Fig. 3d, it appears that the shape of the etched 271 structures reproduces exactly the structural modification gradient.

At longer dissolution times (at t = 1200 s), the structure's height reaches a plateau as shown in Fig. 4. This demonstrates that after this time, the whole poled layer thickness of the C-Na_{5.5} sample was dissolved, and no further contrast between IC and NIC zones can be achieved. On this graph is also shown the maximal height measured by AFM as a function of the dissolution times for the two other selected samples, namely A-Na₂ and C-Na₂. It appears that for short dissolution times (under 300 s), the generated structures are rather small (about 500 nm) for the A-Na₂ sample, bigger for C-Na₂ (2.1 μ m), and the biggest for C-Na_{5.5} (3.7 μ m). If the later reached a plateau after 300 s, it is not the case of the other two samples. The step height continues increasing for C-Na₂ and A-Na₂, and respectively reaches 3.8 and 1.6 μ m after 1200 s.

As predicted, in order to generate large topological contrasts, two parameters must be taken into account: (a) a significant difference in dissolution speeds of poled and unpoled regions and (b) a thick poled layer. From Fig. 3, it appears that the difference in dissolution speeds plays a major role during short dissolution times (C-Na_{5.5} exhibits the biggest structures) while the thickness of the poled layer is a key parameter at longer dissolution times (C-Na₂ and C-Na_{5.5} exhibit similar structure heights after 1200 s).



Fig. 3. (a) Optical image of the electrode employed for micro-poling with conductive and non-
conductive zones highlighted. (b) Normalized Raman spectra extracted from the regions of the
poled C-Na5.5 glass that were in contact (IC) - in red - and non in contact (NIC) - in black -
with the conductive zones of the electrode during the poling process and the Raman difference
spectrum (spectrum IC minus NIC). (c) Spatial evolution of the selected bands from the
Raman difference spectrum. (d) Intensity profile along the X position of the band (i)'s Raman
intensity. (e) Topology of the sample measured by AFM in the same region upon dissolution
for different times. The dotted lines in the two last graphs correspond to the metallic grid
position during the poling treatment



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Fig. 4: Height of the step measured between the region in contact (IC) with the metallic part of the electrode during micro-poling and the one non in contact (NC) for the selected micro-poled samples as a function of the dissolution time. The dotted lines are guides for the eye, the experimental error is estimated to be of 10 %.

302 4. Discussion

Throughout the present work, we have demonstrated that in the glassy Ge-Sb-S-Na system, micro-poling followed by immersion in an amine solvent yields a structured topology of these glasses. A topological contrast as large as about four microns was obtained and is at least four times bigger than anything obtained by this process on silicate glasses. [22–25] Here we discuss the different parameters associated with the origin of this selective etching to shine the light on the different levers that one can have access to in order to control these structures.

Previous structural studies of this exact system [31,36] or related glassy systems [34,35] relying on Raman spectroscopy have shown that the higher the sulphur content, the more numerous the S-S homopolar bonds were. There are no S-S bonds in the stoichiometric glass (A-Na₂) since only one sulphur atom connects the different glass entities, forming $SbS_{3/2}$ pyramids and $GeS_{4/2}$ tetrahedra. [36] When sulphur is added, the interconnections between the different entities are made with more sulphur atoms. Sulphur atoms also form polymeric S_n chains as well as S₈ rings for high sulphur containing compositions.

316 As mentioned earlier, chalcogenide thin film synthesis via spin coating of dissolved glass 317 solution has been the object of many studies with the objective of fabricating new planar 318 photonic components. Dissolution mechanisms of these glasses in amine solvent were therefore 319 at the center of interest of some of these studies. [39,40] They revealed that S-S homopolar 320 bond are preferentially attacked by amine solvents, whereas heteropolar bonds are less reactive. 321 These findings are also confirmed by studies focusing on sulfur-amine chemistry. [41] It 322 follows then that the dissolution rate of glasses exhibiting a high S-S bonds content would be 323 higher, consistent with what we observe with a slower dissolution rate of the stoichiometric 324 glass (A-Na₂) as compared to that of the sulphur-rich glass (D-Na₂).

325 The comparison of the poled and virgin glasses' dissolution rates (Fig. 2) has revealed that in 326 EDA, the poled zone was dissolved more rapidly than the unmodified one. One must note that 327 this is not always the general behavior; for instance, acidic etching of poled silicate results in a 328 faster dissolution rate of the unmodified region compared to the poled one. [22] In our previous 329 work, the Raman study of these glasses (i) as a function of the sodium content and (ii) before 330 and after poling has shown that the structural organization of a poled glass was similar to that 331 of the glass without sodium. [31] In other words, the glass matrix was flexible enough to 332 rearrange itself under the electric field and temperature combination to exhibit the composition and structure of a glass without sodium. However, the dissolution of virgin glasses with
 different sodium contents (Fig. 1b) has shown that the influence of the sodium content on the
 dissolution rate was negligible. Therefore, the poled region's selective etching cannot be
 explained by an effect of the structure nor composition.

337 A previous work on these glassy systems has demonstrated a control of the index variation by 338 micro-poling. [30] A decrease in refractive index was observed in the poled regions (therefore, 339 in the IC zones of the present work). Furthermore, on other glass systems, [33] we have 340 demonstrated that this decrease in refractive index was linked to a decrease in the glass network's density upon poling. The higher the initial sodium content was, the larger the density 341 342 decrease was. Recall, the difference in dissolution rates of poled and unmodified regions is 343 larger for the sample C-Na_{5.5} than for C-Na₂; that is, for a same S/Ge ratio, the difference is 344 larger for higher initial sodium content. It follows that the selective etching of the poled region 345 could be explained by this decrease in density that would facilitate the dissolution

One of the major advantages of this technique is illustrated in Fig. 5a: large areas can be structured at once. So far, only millimeter square surfaces were tested but as thermal poling can be readily scaled up to larger sizes in an efficient manner, it should also be extendable toward centimeter square areas. In Fig. 5b and Fig. 5c, one can see different kind of shapes that can be obtained by micro-poling and subsequent dissolution. This clearly demonstrates the versatility of this process with an accurate spatial control of the structures shapes that can be tuned easily by a control of the immersion time.

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Fig. 5: (a) (top) Topography of a glass sample after micro-poling and dissolution measured by an optical profilometer giving evidence of the possibility to control the topology of large surfaces and (bottom) profile extracted from this image showing the regularity of the process. (b) AFM measured topographies of poled glasses at different dissolution times illustrating the different surface morphologies that can be obtained by this process. (c) Optical image (left) and topography (center and right) measured by an optical profilometer of a poled glass after dissolution. The structured electrode employed here consisted of lithographically made gold spiral. As opposed to previous patterns, here the conductive part is small (only the spiral). This kind of designs open possibilities in the domain of microfluidics for instance.

365 5. Conclusion

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366 Throughout this work, we demonstrate a control of the topology of chalcogenide glasses by the 367 combination of an electrothermal process and a chemical attack. The origin of the selective 368 etching of the poled glasses was discussed: a decrease in the density upon poling is most likely 369 at its origin. The influence of the composition (sulphur and sodium contents) was also 370 evidenced. Significant (up to 4 µm) topological contrasts were achieved, and large (millimeter 371 square) areas were structured. We have shown than a control of the spatial variations in size 372 and shape (in 3-dimensions) was possible by tailoring the glass type, the electrode feature and 373 the experimental conditions (of the poling and of the chemical attack). This process could be 374 seen as an alternative to classical lithography techniques.

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386 Disclosures

387 The authors declare no conflict of interest.

388 Data availability

Bata underlying the results presented in this paper are not publicly available at this time but may be obtained from the authors upon reasonable request.

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