

New Two-Photon Absorbing Squaraine Derivative with Efficient Near-Infrared Fluorescence, Superluminescence, and High Photostability

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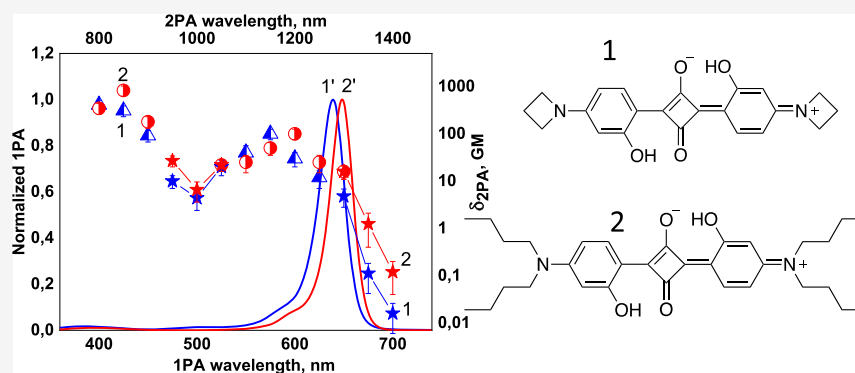


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ABSTRACT: The nature of linear photophysical and nonlinear optical properties of a new squaraine derivative 2,4-bis[4-(azetidyl)-2-hydroxyphenyl]squaraine (**1**) with efficient near-infrared (NIR) emission was comprehensively analyzed based on spectroscopic, photochemical, and two-photon absorption (2PA) measurements, along with quantum chemical analysis. The steady-state absorption, fluorescence, and excitation anisotropy spectra of **1** and its fluorescence emission lifetimes revealed the multiple aspects of the electronic structure of **1**, including the relative orientations of the main transition dipoles, effective rotational volumes in solvents of different polarities, and a maximum molar extinction of $1.35 \times 10^5 \text{ M}^{-1} \cdot \text{cm}^{-1}$, which is unusually small for similar symmetric squaraines. The degenerate 2PA spectrum of **1** was obtained over a broad spectral range under femtosecond excitation, using standard open-aperture Z-scan and two-photon induced fluorescence methods, revealing maximum 2PA cross sections of $\sim 400 \text{ GM}$. Squaraine **1** exhibited efficient superluminescence emission in the polar solvent (dichloromethane) at room temperature under femtosecond pumping conditions. Quantum chemical analysis of the electronic structure of **1** was performed using the DFT/TD-DFT level of theory and found to be in good agreement with experimental data. The new squaraine derivative **1** displayed high fluorescence quantum yield, efficient NIR superluminescence, large 2PA cross sections, and high photostability with a photodecomposition quantum yield $\sim 4 \times 10^{-6}$, suggesting its potential for applications in two-photon fluorescent bioimaging and lasing.

1. INTRODUCTION

The development of new squaraine derivatives with intense linear absorption bands, efficient near-infrared (NIR) emission, large two-photon absorption (2PA) cross sections, and high photostability is a subject of interest for a number of important scientific and technological areas, including nonlinear optics,¹ optical data recording,² optoelectronics,^{3–5} metal-ion sensing,⁶ photodynamic therapy,^{7,8} and fluorescence bioimaging,^{9–11} among others. In general, squaraine derivatives can be presented as resonantly stabilized zwitterionic structures^{8,12,13} with an electron-deficient central squarainyl ring C_4O_2 and electron-donating termini.^{14–16} Certain squaraine molecules exhibit high photochemical stability^{17–19} and possess a wide

variety of specific linear and nonlinear spectroscopic properties determined by the specific structure of the molecular architecture, including intramolecular steric factors,^{4,20,21} hydrogen bonding,²² symmetry of the electronic distribution,^{19,23,24} and parameters of the terminal groups.^{13,16,25–27}

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The majority of symmetrical squaraine derivatives (with corresponding quadrupolar molecular structure: donor–acceptor–donor^{15,28}) are characterized by relatively sharp, intense long-wavelength absorption bands in the red-NIR spectral range (maximum extinction coefficient, $\epsilon^{\max} \sim (3\text{--}4) \times 10^5 \text{ M}^{-1} \text{ cm}^{-1}$),^{18,29,30} high fluorescence quantum yields, $\Phi_{\text{f}} > 0.6$,^{18,31,32} and some of the highest single-molecule 2PA efficiencies (corresponding 2PA cross sections, $\delta_{2\text{PA}}^{\max} \sim 10^4 \text{ GM}$).^{11,33–35} The nature of these large values of $\delta_{2\text{PA}}^{\max}$ for symmetrical squaraines is mainly determined by the large $\epsilon^{\max}$ and the possibility to realize double-resonance excitation conditions.³⁵ Squaraine derivatives with unsymmetrical molecular structures often exhibit much broader absorption bands with smaller values of $\epsilon^{\max}$ ³⁶ and possess significant potential for applications in organic solar cell engineering.^{5,23,37}

All types of squaraine molecules typically reveal a strong decrease in emission intensity with the increase in polarity of the surrounding medium^{20,26,30–33,38} that can be explained by the increased efficiency of the nonradiative vibronic relaxations, which dramatically reduce their fluorescence. In this study, we present the synthesis and comprehensive characterization of this new 2,4-bis[4-(azetidyl)-2-hydroxyphenyl]squaraine (**1**) with potential for two-photon fluorescence microscopy applications.^{9,11,17} Linear photophysical, photochemical, 2PA, and superluminescence properties of **1** were investigated in comparison with a known analogous squaraine structure 2,4-bis[4-(*N,N*-dibutylamino)-2-hydroxyphenyl]squaraine³⁹ (**2**) in liquid media at room temperature (see structures in Figure 1). The electronic

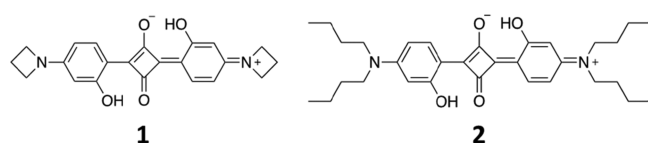


Figure 1. Molecular structure of squaraine dyes.

properties of **1** and **2** were also analyzed with density functional theory (DFT)/time-dependent DFT (TD-DFT) quantum-chemical calculations, and the essential excited states, oscillator strengths of the corresponding electronic transitions, main molecular orbital configurations, and the role of the end substituents are described.

2. EXPERIMENTAL SECTION

2.1. Synthesis of Azetidyl-Functionalized and Reference Squaraines. The azetidyl starting material and reference squaraine **2** were prepared using previously reported methods,^{40,41} and the ¹H and ¹³C NMR results were consistent with the reported data. The main synthetic procedures for **1** and **2** are presented in Schemes 1 and 2, respectively.

2.1.1. Synthesis of 2,4-bis[4-(azetidyl)-2-hydroxyphenyl]squaraine (1**).** 3-(Azetid-1-yl)phenol was prepared using a previously reported method.⁴⁰ A round-bottom flask was

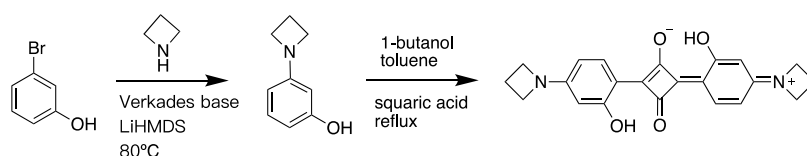
charged with Pd(OAc)₂ (0.579 mmol) in toluene (30 mL) and backfilled with nitrogen. Sequentially, 3-bromophenol (11.6 mmol), 2,8,9-triisobutyl-2,5,8,9-tetraaza-1-phosphabicyclo[3.3.3]undecane (“Verkade base”, 1.16 mmol), and LiHMDS (1.0 M in THF, 26.6 mmol) were added to the flask, after which azetidine (13.9 mmol) was added slowly, and the reactants were stirred at 80 °C for 18 h. The mixture was cooled to room temperature, deposited onto silica, and washed with hexane and ethyl acetate. The solvent was removed under pressure and purified by silica gel chromatography (70:30, hexanes:ethyl acetate). ¹H NMR of 3-(azetid-1-yl)phenol (500 MHz, CDCl₃) δ : 7.04 (t, 1H), 6.19 (ddd, 1H), 6.05 (ddd, 1H), 5.92 (t, 1H), 4.77 (s, 1H), 3.85 (t, 4H), 2.35 ppm (p, 2H). ¹³C NMR of 3-(azetid-1-yl)phenol (500 MHz, CDCl₃) δ : 156.48, 153.78, 129.94, 104.54, 104.27, 98.49, 52.40, 16.86 ppm.

The final product **1** was prepared by charging a round-bottom flask, fitted with a Dean–Stark apparatus, with squaric acid (1 mol) in a 1-butanol (15 mL) and toluene (6 mL) mixture. The solution was heated to 100 °C for 15 min to allow the squaric acid to dissolve. Subsequently, 3-(azetid-1-yl)phenol (2 mol) in toluene (8 mL) was added dropwise over 20 min. The solution was heated to reflux for 2 h. Blue solid was obtained through filtration and washed with ethyl acetate. The product was extracted from the solid by washing with chloroform. ¹H NMR of **1** (500 MHz, CDCl₃) δ : 12.17 (s, 2H), 7.89 (d, 2H), 5.99 (dd, 2H), 5.77 (d, 2H), 4.15 (t, 8H), 2.48 ppm (quintet, 4H). HR-MS (ESI) (*m/z*): [*M* + *H*] + Calculated for C₂₂H₂₁N₂O₄, 377.1496; Found, 377.1492.

2.1.2. Synthesis of 2,4-bis[4-(*N,N*-dibutylamino)-2-hydroxyphenyl]squaraine (2**).** A round-bottom flask, fitted with a Dean–Stark apparatus, was charged with squaric acid (1 mol) in a 1-butanol (15 mL) and toluene (6 mL) mixture. The solution was heated to 100 °C for 15 min to allow the squaric acid to dissolve. Subsequently, *N,N*-(dibutylamino)phenol (2 mol) in toluene (8 mL) was added dropwise over 20 min. The solution was heated to reflux for 2 h. Green crystals were obtained through filtration and washed with ethyl acetate. ¹H NMR of **2** (500 MHz, CDCl₃) δ : 12.15 (s, 2H), 7.92 (d, 2H), 6.37 (dd, 2H), 6.17 (d, 2H), 3.43 (t, 8H), 1.68 (quintet, 8H), 1.44 (sextet, 8H), 1.03 ppm (t, 12H). ¹³C NMR of **2** (500 MHz, CDCl₃) δ : 183.07, 172.17, 164.13, 156.40, 132.85, 109.83, 107.61, 98.53, 51.41, 29.85, 20.26, 13.93 ppm.

2.2. Linear Photophysical and Photochemical Characterization of **1 and **2**.** All linear photophysical and photochemical measurements were performed in air-saturated spectroscopic-grade dichloromethane (DCM) and toluene (TOL), which were purchased from commercial sources and used without further purification. The steady-state linear or one-photon absorption (1PA) spectra of **1** and **2** were recorded with a high-performance Cary 5000 UV–vis–NIR (Agilent) spectrophotometer using 1 cm path length quartz cuvettes and squaraine concentrations, *C* $\approx (3\text{--}5) \times 10^{-6} \text{ M}$. The steady-state and time-resolved fluorescence measure-

Scheme 1. Synthetic Route for Azetidyl-Functionalized Squaraine **1**



Scheme 2. Synthetic Route for Reference Squaraine 2

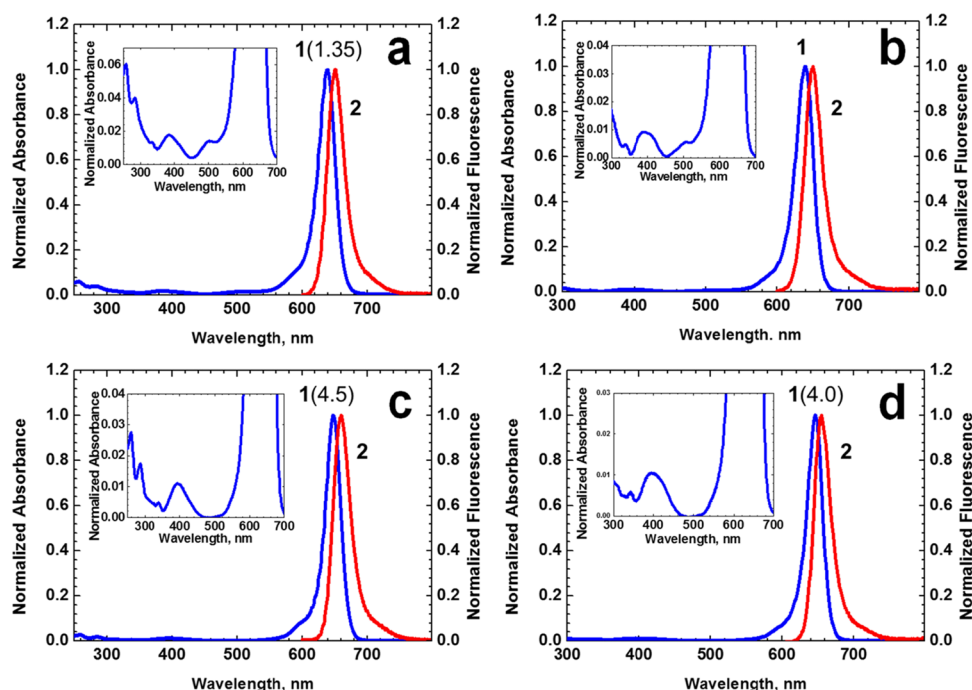
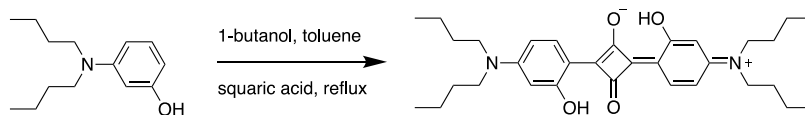


Figure 2. Normalized steady-state 1PA (1) and fluorescence (2) spectra of **1** in DCM (a), TOL (b), **2** in DCM (c), and TOL (d). Number in the parentheses represents the maximum extinction coefficient ($\epsilon^{\text{max}} \times 10^{-5}$, $\text{M}^{-1}\cdot\text{cm}^{-1}$).

ments, including corrected fluorescence and excitation anisotropy spectra values of the fluorescence quantum yields, Φ_{fl} , and emission lifetimes, τ_{fl} , were determined using a FLS980 spectrofluorimeter (Edinburg Instruments Ltd.) with dilute solutions ($C \approx (1-2) \times 10^{-7}$ M) and standard spectrofluorimetric quartz cuvettes. Fluorescence quantum yields of **1** and **2** were determined using a standard comparative method⁴² with cresyl violet in methanol as a reference ($\Phi_{\text{fl}} \approx 0.54$).⁴³

The fundamental excitation anisotropy spectra, $r_0(\lambda)$, for **1** and **2** were obtained at room temperature in viscous polytetrahydrofuran (pTHF), where the molecular rotational correlation time, $\theta = \eta \cdot V/kT$ (η , V , k , and T are the solvent viscosity, effective rotational volume of the molecule, Boltzmann constant, and absolute temperature, correspondingly)⁴⁴ was much larger than τ_{fl} and rotational depolarization effects were negligible. In this case, the experimentally observed excitation anisotropy spectrum $r(\lambda) = r_0(\lambda)/(1 + \tau_{\text{fl}}/\theta) \approx r_0(\lambda)$.⁴² Photochemical stabilities of **1** and **2** were investigated in air-saturated TOL and DCM under continuous wave (CW) laser irradiation at $\lambda = 637$ nm with average beam irradiance ≈ 500 mW/cm². Corresponding values of the photodecomposition quantum yields, Φ_{ph} (i.e., the number of decomposed molecules divided by the number of absorbed photons), were determined by a known absorption method⁴⁵ as

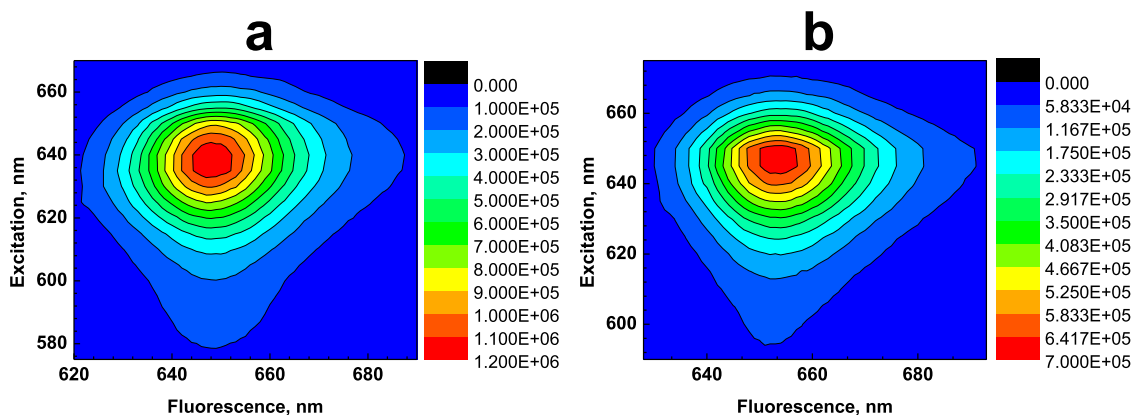
$$\Phi_{\text{ph}} = \frac{[D(\lambda, 0) - D(\lambda, t_{\text{ir}})] \cdot N_A}{\epsilon(\lambda) \cdot 10^3 \cdot \int_{\lambda} \int_0^{t_{\text{ir}}} I_0(\lambda) \cdot [1 - 10^{-D(\lambda, t)}] \cdot d\lambda \cdot dt} \quad (1)$$

where $D(\lambda, 0)$, $D(\lambda, t_{\text{ir}})$, $\epsilon(\lambda)$, N_A , λ , t_{ir} , and $I_0(\lambda)$ are the values of initial and final absorbance of the sample solution, extinction coefficient (in $\text{M}^{-1}\cdot\text{cm}^{-1}$), Avogadro's number, excitation wavelength, irradiation time, and spectral distribution of the excitation irradiance, respectively.

2.3. 2PA Spectra and Superluminescence Measurements. Nonlinear optical properties of **1** and **2** were investigated in DCM at room temperature using femtosecond laser systems and corresponding experimental techniques, described in detail previously.^{17,19,46} The degenerate 2PA spectra were obtained over a broad spectral range (800–1400 nm) by open-aperture Z-Scan⁴⁷ and two-photon-induced fluorescence (2PF) methods⁴⁸ using 1 and 10 mm path length quartz cuvettes with dye concentrations $C \sim 10^{-3}$ and $C \sim 10^{-5}$ M, respectively. The exit beam from the optical parametric amplifier (OPA, TOPAS C, Light Conversion Ltd., with pulse energy, $E_p \leq 250$ μJ , pulse duration, $\tau_p \approx 150$ fs, and repetition rate 1 kHz), which was pumped by the regenerative amplifier (Clark MXR, CPA 2010), was employed for Z-scan measurements. The same laser beam in combination with the PTI spectrofluorimeter was used for 2PF measurements. The superluminescence of **1** and **2** was observed from the 10 mm spectrofluorimetric quartz cuvettes with dye concentrations $\sim 10^{-4}$ M under femtosecond pumping by the second harmonic of the OPA tuned to 1300 nm. The pump laser beam (excitation wavelength, $\lambda_{\text{exc}} \approx 650$ nm, $E_p \leq 5$ μJ)

Table 1. Linear Spectroscopic and Photochemical Parameters of New Squaraine 1 and Reference One 2 in TOL and DCM^a

compound	1		2	
solvent	TOL	DCM	TOL	DCM
λ_{ab}^{max} , nm	639 $\pm$ 1	639 $\pm$ 1	646 $\pm$ 1	648 $\pm$ 1
λ_{fl}^{max} , nm	650 $\pm$ 1	651 $\pm$ 1	654 $\pm$ 1	660 $\pm$ 1
Stokes shift, nm (cm^{-1})	11 $\pm$ 2 (260 $\pm$ 30)	12 $\pm$ 2 (290 $\pm$ 30)	8 $\pm$ 2 (190 $\pm$ 20)	12 $\pm$ 2 (280 $\pm$ 20)
$\epsilon^{max} \times 10^{-5}$, M ⁻¹ ·cm ⁻¹		1.35 $\pm$ 0.15	4.5 $\pm$ 0.3	4.0 $\pm$ 0.2
μ_{01} , Debye		8.8	14.2	14.0
Φ_{fl}	0.97 $\pm$ 0.05	0.77 $\pm$ 0.05	0.88 $\pm$ 0.05	0.86 $\pm$ 0.05
τ_{fl} , ns	2.5 $\pm$ 0.2	2.6 $\pm$ 0.2	2.5 $\pm$ 0.1	2.6 $\pm$ 0.2
τ_{fl}^{cal} , ns		4.8 $\pm$ 1	1.6 $\pm$ 0.3	2.1 $\pm$ 0.4
$r(\lambda_{ab}^{max})$	0.016 $\pm$ 0.003	0.015 $\pm$ 0.003	0.026 $\pm$ 0.005	0.020 $\pm$ 0.004
θ , ns	0.15 $\pm$ 0.05	0.14 $\pm$ 0.04	0.21 $\pm$ 0.06	0.17 $\pm$ 0.05
V , Å ³	1050	1450	1470	1760
$\Phi_{ph} \times 10^6$		3.8 $\pm$ 0.8	0.02 $\pm$ 0.004	1.4 $\pm$ 0.3

**Figure 3.** 3D fluorescence emission maps of 1 (a) and 2 (b) in TOL.

was focused by a cylindrical lens to a waist of 0.15×10 mm while superluminescence emission was detected in the transverse direction by a fiber optic spectrometer (HR4000, Ocean Optics, Inc.).

2.4. Computational Analysis. The nature of the electronic properties of 1 and 2 was estimated by quantum-chemical calculations, which were performed with the Gaussian 2009, Rev.A02.⁴⁹ DFT with the 6-31G(d,p) basis set and the B3LYP functional was used for geometry optimization. TD-DFT was employed to describe the excited state (including the optimized geometry). The polarizable continuum model (PCM)⁵⁰ has been used to take into account the surrounding solvent. Based on the known negligible effect of alkyl substituents on the electronic properties of symmetrical squaraines,⁵¹ the *n*-butyl groups on the nitrogens in model compound 2 were replaced by methyl groups. All linear spectral parameters were predicted using optimized ground and first excited singlet electronic-state molecular geometry for absorption and emission spectra, respectively.

3. RESULTS AND DISCUSSION

3.1. Linear Spectral and Photochemical Properties of 1 and 2. The main photophysical characteristics and photodecomposition quantum yields, Φ_{ph} , of the new squaraine derivative 1 and reference compound 2 are presented in Figure 2 and Table 1. The linear 1PA spectra of 1 and 2 exhibit relatively sharp (FWHM $\sim$ 30 nm) and intense ($\epsilon^{max} > 1.3 \cdot 10^5$ M⁻¹ cm⁻¹) long-wavelength absorption bands with maxima at $\sim$ 640–680 nm (Figure 2, curves 1),

which are typical for symmetrical squaraines in this spectral range.^{15,18,20} Very weak short-wavelength linear absorption bands at $\sim$ 400 nm were observed for 1 and 2 in both solvents studied (see insets in Figure 2), which potentially can be active in 2PA processes.^{15,35} It should be mentioned that a large difference in the values of ϵ^{max} was obtained for 1 and 2 in DCM (Table 1 showing ϵ^{max} (1) < ϵ^{max} (2)), which is not clearly understood because of the similarity of the molecular structures and all other linear photophysical parameters determined in this work. The solubility of 1 in TOL was relatively low compared to 1 and 2 in DCM while no evidence of aggregation was observed. Hence, these processes cannot be a reason for the large difference in ϵ^{max} . Fluorescence emission spectra of squaraine 1 were similar to those of reference compound 2 in TOL and DCM with small Stokes shifts (<15 nm) and full independence of the excitation wavelength (i.e., without violation of Kasha's rule⁴²). Three-dimensional (3D) fluorescence maps for 1 and 2 (Figure 3) confirm this assertion along with the full independence of the corresponding excitation spectra of the observed wavelength. The fluorescence emission quantum yields, Φ_{fl} , were relatively high (see Table 1), especially for 1 in TOL ($\Phi_{fl} \approx 1.0$), which can be employed as a fluorescence reference standard in the red-NIR spectral range. The kinetics of fluorescence decay of 1 and 2 was characterized by a single-exponential process in TOL and DCM (Figure 4) with the characteristic lifetimes $\tau_{fl} \approx 2.5$ –2.6 ns, which were nearly the same for both squaraines and does not explain the nature of such a large difference in their ϵ^{max} and corresponding transition dipoles,

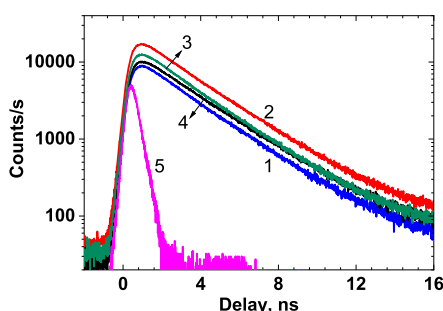


Figure 4. Kinetics of fluorescence decay for squaraine **1** (1, 2) and **2** (3, 4) in TOL (1, 3) and DCM (2, 4), and instrument response function (5).

$\mu_{01} \approx 0.096 \cdot \sqrt{\int \epsilon(\nu) \cdot \lambda_{ab}^{max} \cdot d\nu}$, estimated from the experimental data (μ_{01} in Debye, extinction coefficient in the main absorption band, $\epsilon(\nu)$ in $M^{-1} \cdot cm^{-1}$, λ in cm, and ν in cm^{-1}).⁵²

^aSteady-state absorption λ_{ab}^{max} and fluorescence λ_{fl}^{max} maxima; values of Stokes shifts; maximum extinction coefficients, ϵ^{max} ; fluorescence emission quantum yields, Φ_{fl} ; estimated transition dipoles for $S_0 \rightarrow S_1$ electronic transition, μ_{01} (S_0 and S_1 are the ground and first excited electronic state, correspondingly);⁵² experimentally observed, τ_{fl} , and calculated, τ_{fl}^{cal} , fluorescence lifetimes;⁵³ excitation anisotropies, $r(\lambda_{ab}^{max})$, rotational correlation times, θ , effective rotational volume, V , and photodecomposition quantum yields, Φ_{ph} .

Experimentally observed fluorescence lifetimes τ_{fl} can be also calculated from the linear spectral data of **1** and **2** as⁴² $\tau_{fl}^{cal} = \tau_N \cdot \Phi_{fl}$, where τ_N is the natural lifetime that can be determined from the Strickler–Berg theoretical approach.⁵³ The values of τ_{fl}^{cal} are in a good agreement with experimental ones for reference squaraine **2** (see Table 1), which is evidence of relatively small changes in the molecular geometry under electronic excitation $S_0 \rightarrow S_1$ and fulfillment of the mirror symmetry between absorption and emission contours.⁵⁴ In contrast, a large discrepancy was observed for the corresponding lifetimes for **1**, which can be related to an unusually small (for symmetrical squaraines) value of ϵ^{max} . The steady-state fundamental excitation anisotropy spectra of **1** and **2**, $r_0(\lambda)$, were obtained in viscous pTHF (Figure 5, curves 1) that reveal the nature of the main 1PA band: nearly constant values of $r_0(\lambda) \approx const$ in the spectral range ≈ 560 – 670 nm are related to the dominant role of a single electronic transition $S_0 \rightarrow S_1$ in the main long-wavelength absorption band of **1** and **2**.⁵⁵ The maxima values of $r_0(\lambda) \approx 0.3$ provide an estimation of the

angle between the absorption, μ_{01} , and emission, μ_{10} , transition dipoles as $\approx 25^\circ$, which is too large to allow for a correct determination of the relative orientations of the corresponding absorption dipoles μ_{01} and μ_{0n} (n is the number of higher excited electronic state).⁵⁵ Excitation anisotropy spectra $r(\lambda)$ were also determined for **1** and **2** in low-viscosity TOL and DCM (Figure 5, curves 2, 3) in order to determine $r(\lambda_{ab}^{max})$, rotational correlation times, θ , and effective rotational volume, V (see Section 2.2 and Table 1). It should be noted that within the main absorption band, the signal to noise in the low-viscosity solvents is significant. According to these data, the values of θ exhibit a relatively weak dependence on solvent polarity in contrast to the effective rotational volumes, V , which exhibit a noticeable increase in the size of the solvate cage in more polar solvents for both of the investigated squaraines.

The level of the photochemical stability of **1** and **2** was investigated quantitatively in air-saturated TOL and DCM at room temperature using the absorption method⁴⁵ with low-intensity laser irradiation in the main absorption band at $\lambda \approx 637$ nm. Accordingly, small changes in the main absorption bands of **1** and **2** (Figure 6) were detected during the irradiation process, and the corresponding values of the photodecomposition quantum yields, Φ_{ph} , were calculated by eq 1 and shown in Table 1. It should be mentioned that the kinetics of the observed photodecomposition processes corresponded to a first-order photoreaction,⁵⁶ and no substantial photoproducts were detected near the 1PA maxima of **1** and **2**. From these data, the highest photostability was observed for **2** in TOL ($\Phi_{ph} \sim 10^{-8}$) and dramatically decreased up to two orders of magnitude in polar DCM. The level of photostability of new squaraine **1** is somewhat lower than that of **2** but still comparable with the best laser dyes in liquid media⁵⁷ and can be used for important practical applications, including as a contrast agent in laser scanning fluorescence microscopy techniques.⁹

3.2. 2PA Properties and Superluminescence Effects in Liquid Solutions of 1 and 2. The values of 2PA cross sections, σ_{2PA}^{max} , were obtained for **1** and **2** in DCM using direct open-aperture Z-Scan⁴⁷ and relative 2PF⁴⁸ methods, while the corresponding degenerate spectra are shown in Figure 7. As follows from these data, the shape of the obtained 2PA spectra is similar to some known symmetric squaraines^{18,32,34} with the weak maximum ($\delta_{2PA}^{max} \approx 90$ – 100 GM) at ~ 1150 to 1200 nm and strong increase (up to 400 – 830 GM) in the short-wavelength spectral range at ~ 800 – 850 nm. It is worth noting

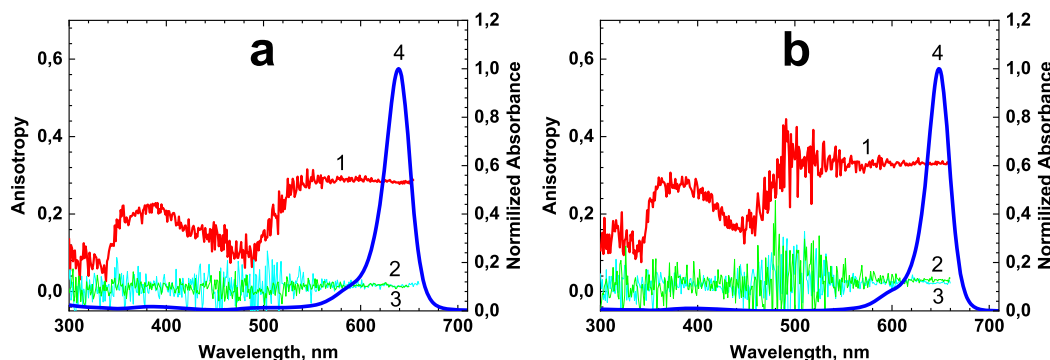


Figure 5. Steady-state excitation anisotropy spectra of squaraine **1** (a) and **2** (b) in pTHF (1), TOL (2), and DCM (3); corresponding linear 1PA spectra in DCM (4).

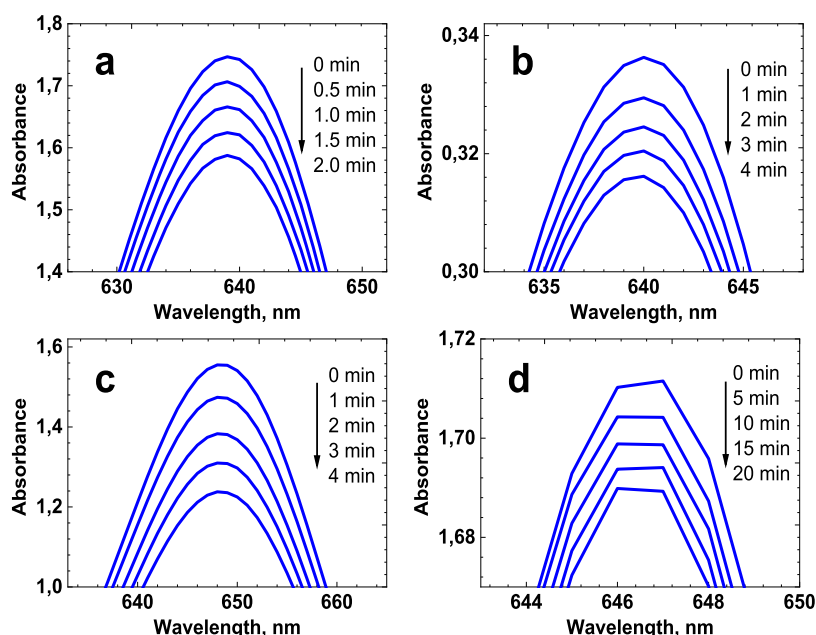


Figure 6. Photobleaching of the main linear absorption band of **1** (a,b) and **2** (c,d) in DCM (a,c) and TOL (b,d) under CW laser irradiation at 637 nm with beam irradiance $\approx 500 \text{ mW/cm}^2$. Corresponding irradiation times, t_{ir} , are indicated on the right.

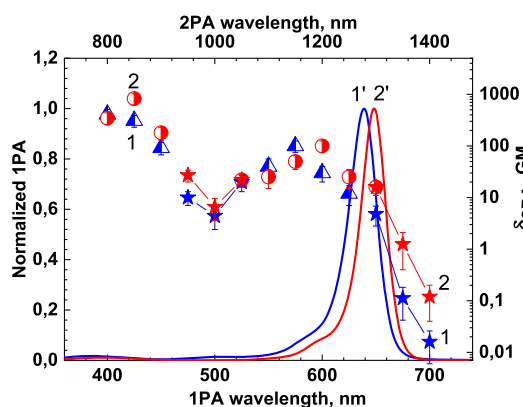


Figure 7. Degenerate 2PA (**1**, **2**) and normalized linear 1PA spectra (**1'**, **2'**, solid lines) of **1** (**1'**) and **2** (**2'**) in DCM. Circles are for **2** using Z-scan and triangles are for **1** using Z-scan. Stars show data (**1** blue, **2** red) obtained using the 2PF method. The 2PF data at 1050 nm and 1300 nm are calibrated using the Z-scan data at these wavelengths.

that two-photon transitions in the main linear absorption band are strictly forbidden for centrosymmetric molecules.⁵⁸ Nevertheless, the spectral position of the weak 2PA band roughly corresponded to the vibrational shoulder in the linear 1PA spectrum at ~ 580 to 600 nm (Figure 7, curves **1'**, **2'**), and the nature of this band can be assigned to vibronic interactions with the asymmetric vibrational modes of **1** and **2**.^{15,59} The shape of excitation anisotropy spectra (Figure 5, curves **1**) revealed some higher excited electronic levels, S_n , in the spectral range $380\text{--}420 \text{ nm}$, which are in close proximity to the “double-resonance” excitation conditions for $S_0 \rightarrow S_n$ two-photon transitions,³⁵ and can explain a noticeable increase in the 2PA efficiency in this spectral range. In general, new squaraine **1** has 2PA properties similar to reference **2** and, according to a “figure of merit”,¹⁸ ($F_M = \delta_{2PA} \cdot \Phi_R / \Phi_{ph} \approx 10^8$), exhibits high potential for application in two-photon fluorescence bioimaging.

The potential abilities of the squaraines **1** and **2** for light amplification were investigated in liquid solutions at room temperature under one-photon femtosecond transverse pumping in the main 1PA band with pump pulse energy $E_p \leq 5 \mu\text{J}$ (see Section 2.3). Efficient superluminescence emission was observed from relatively concentrated DCM solutions ($C \sim 10^{-4} \text{ M}$). Typical spectra are shown in Figure 8. The spontaneous fluorescence contours (Figure 8a,c, curves **4**) were strongly reabsorbed under the employed concentrations and, consequently, transformed into much narrower superluminescence emission bands (curves **1**, FWHM $< 15 \text{ nm}$) with an increase in E_p . According to the data, the dependence of the collected emission intensity on the pump pulse energy revealed an obvious threshold character (Figure 8b,d) with the threshold values, $E_{th} \approx 0.12\text{--}0.15 \mu\text{J}$, and exhibited linearity in the energy range $E_p < E_{th}$ (see insets in Figure 8b,d). It is worth mentioning that molecular superluminescence emission in bioimaging is a promising approach to expand modern laser scanning fluorescence microscopy methods.^{18,19}

3.3. Theoretical Analysis of the Electronic Properties of **1 and **2**.** The nature of the electronic transitions in **1** and **2** was analyzed based on the DFT/TD-DFT theoretical approach using the Gaussian-2009 Rev.A02 suite of programs⁴⁹ and the PCM model⁵⁰ (see Section 2.4). Optimized molecular geometries in S_0 and S_1 electronic states, corresponding changes in the bond lengths under electronic excitation $S_0 \rightarrow S_1$, and the main calculated spectroscopic parameters of **1** and **2** in DCM are presented in Figure 9 and Table 2, respectively. All optimized geometries in Figure 9 revealed nearly planar molecular structures and the possibility of intramolecular hydrogen bonding between the hydroxyl groups and the oxygens of the squaric center. This is consistent with the experimentally observed high fluorescence quantum yields of **1** and **2** in DCM (see Table 1). The changes in the bond lengths under $S_0 \rightarrow S_1$ excitation were relatively small ($\leq 0.013 \text{ \AA}$) and nearly the same for **1** and **2**. Also, as deduced from the obtained data, the values of the calculated oscillator strengths and transition dipoles exhibit negligible dependence

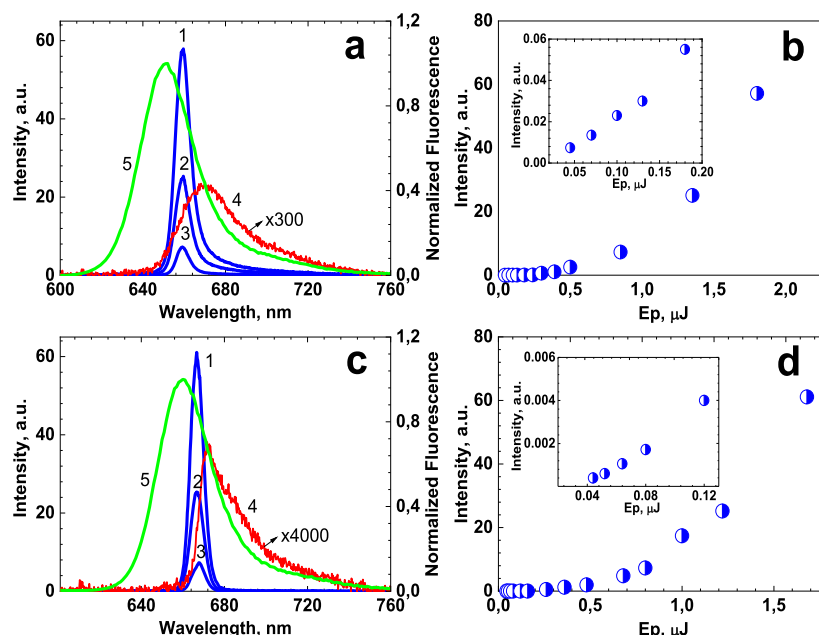


Figure 8. (a) Superluminescence bands of **1** in DCM at $C \approx 2 \cdot 10^{-4}$ M under femtosecond pumping with pulse energy $E_p \approx 1.8 \mu\text{J}$ (1), $1.35 \mu\text{J}$ (2), $0.85 \mu\text{J}$ (3), $0.24 \mu\text{J}$ (4), and normalized steady-state fluorescence spectrum of **1** in DCM at $C \approx 2 \cdot 10^{-7}$ M (5). (b) Dependence of collected emission intensity, I , on E_p for **1** in DCM at $C \approx 2 \cdot 10^{-4}$ M, and the corresponding initial part of this dependence (inset in (b)). (c) Superluminescence bands of **2** in DCM at $C \approx 3 \cdot 10^{-4}$ M under femtosecond pumping with $E_p \approx 1.68 \mu\text{J}$ (1), $1.22 \mu\text{J}$ (2), $0.8 \mu\text{J}$ (3), $0.16 \mu\text{J}$ (4), and normalized steady-state fluorescence spectrum of **2** in DCM at $C \approx 10^{-7}$ M (5). (d) Dependence $I = f(E_p)$ for **2** in DCM at $C \approx 3 \cdot 10^{-4}$ M, and the corresponding initial part of this dependence (inset in (d)).

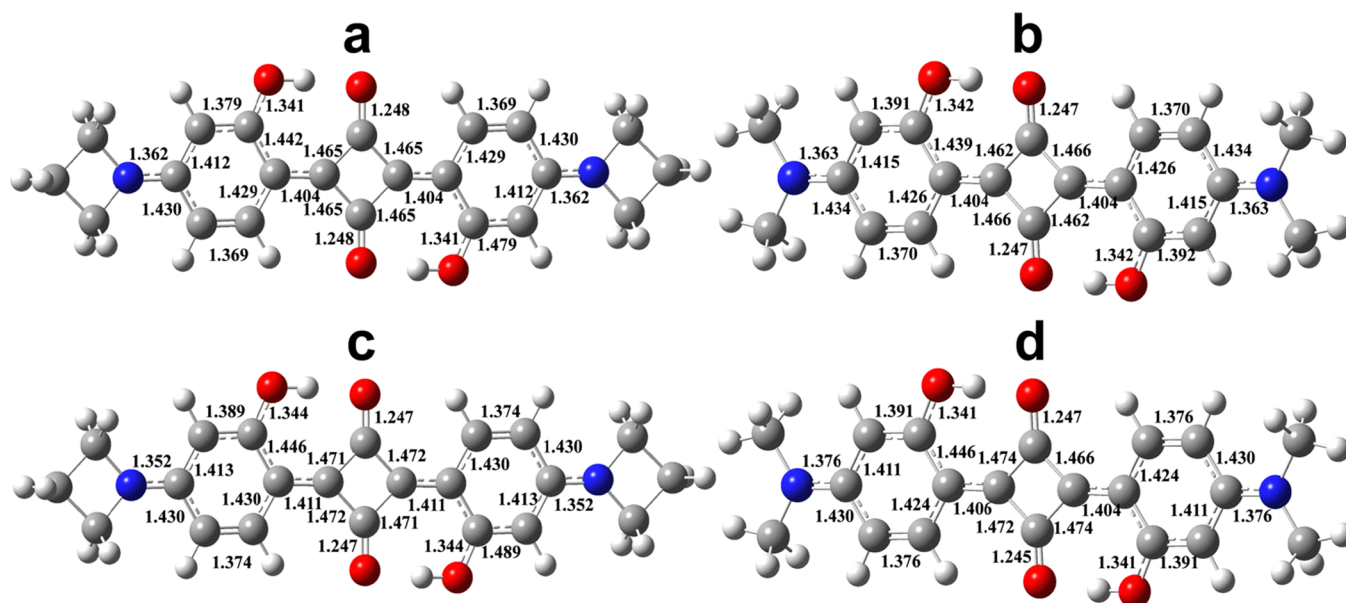


Figure 9. Optimized molecular geometries of **1** (a,c) and **2** (b,d) in DCM for S_0 (a,b) and S_1 (c,d) electronic states.

on the nature of the end substituents and do not explain the large difference in ϵ^{max} experimentally observed for **1** and **2** in DCM (i.e., $\epsilon^{\text{max}}(\mathbf{1}) < \epsilon^{\text{max}}(\mathbf{2})$). The absolute calculated value of the transition dipole μ_{01} for **2** in DCM is in a good agreement with the corresponding one estimated from the experimental spectral data (see Table 1). At the same time, even though solvent effects have been included, there is no perfect match of the calculated electronic transition wavelengths and experimental values of the absorption and emission maxima. Typical for the standard TD-DFT approach employed for cyanine dyes,⁶⁰ nevertheless, this method allows one to

correctly analyze the nature of the first and several higher excited electronic states. Through comparison of the data in Figure 2 and Table 2, the main long-wavelength absorption bands of **1** and **2** correspond solely to the $S_0 \rightarrow S_1$ ($\pi \rightarrow \pi^*$) transition with nearly pure HOMO $\rightarrow$ LUMO configuration (HOMO and LUMO are the highest occupied and the lowest unoccupied molecular orbitals, respectively) and relatively high intensity (oscillator strength, $f \approx 1.75\text{--}1.94$). This is consistent with the experimentally observed constant values of anisotropy in the main IPA bands of **1** and **2** (Figure 5, curves 1) and the large value of $\epsilon^{\text{max}} \approx 4 \times 10^5 \text{ M}^{-1} \text{ cm}^{-1}$ for reference squaraine

Table 2. Calculated Parameters of the Essential Excited States of **1** and **2**: Wavelengths, λ , Oscillator Strength, f , Transition Dipoles μ_{01} , and Leading Configurations (HOMOs and LUMOs Are the Highest Occupied and the Lowest Unoccupied Molecular Orbitals, Respectively)

N/N	transition	λZ , nm	f	μ_{01} , Debye	type	main configurations
absorption of 1	$S_0 \rightarrow S_1$	549.7	1.9402	15.04	$\pi \rightarrow \pi^*$	0.99 HOMO $\rightarrow$ LUMO>
	$S_0 \rightarrow S_2$	403.0	0.0000	0.00	$n \rightarrow \pi^*$	0.98 HOMO-4 $\rightarrow$ LUMO>
	$S_0 \rightarrow S_3$	402.8	0.0768	2.99	$\pi \rightarrow \pi^*$	0.97 HOMO-2 $\rightarrow$ LUMO>
	$S_0 \rightarrow S_4$	394.4	0.0000	0.00	$\pi \rightarrow \pi^*$	0.30 HOMO-3 $\rightarrow$ LUMO> -0.63 HOMO-1 $\rightarrow$ LUMO>
	$S_0 \rightarrow S_5$	350.6	0.0000	0.00	$\pi \rightarrow \pi^*$	0.61 HOMO-3 $\rightarrow$ LUMO> -0.27 HOMO-1 $\rightarrow$ LUMO>
fluorescence of 1	$S_1 \rightarrow S_0$	560.4	1.8998		$\pi^* \rightarrow \pi$	0.71 LUMO $\rightarrow$ HOMO>
absorption of 2	$S_0 \rightarrow S_1$	547.5	1.7882	14.42	$\pi \rightarrow \pi^*$	0.99 HOMO $\rightarrow$ LUMO>
	$S_0 \rightarrow S_2$	411.2	0.0000	0.00	$n \rightarrow \pi^*$	0.98 HOMO-4 $\rightarrow$ LUMO>
	$S_0 \rightarrow S_3$	410.7	0.0769	2.99	$\pi \rightarrow \pi^*$	0.45 HOMO-2 $\rightarrow$ LUMO> +0.49 HOMO-1 $\rightarrow$ LUMO>
	$S_0 \rightarrow S_4$	401.4	0.0000	0.00	$\pi \rightarrow \pi^*$	0.53 HOMO-2 $\rightarrow$ LUMO> -0.42 HOMO-1 $\rightarrow$ LUMO>
	$S_0 \rightarrow S_5$	354.8	0.0000	0.00	$\pi \rightarrow \pi^*$	0.63 HOMO-3 $\rightarrow$ LUMO>
fluorescence of 2	$S_1 \rightarrow S_0$	556.8	1.7550		$\pi^* \rightarrow \pi$	0.71 LUMO $\rightarrow$ HOMO>

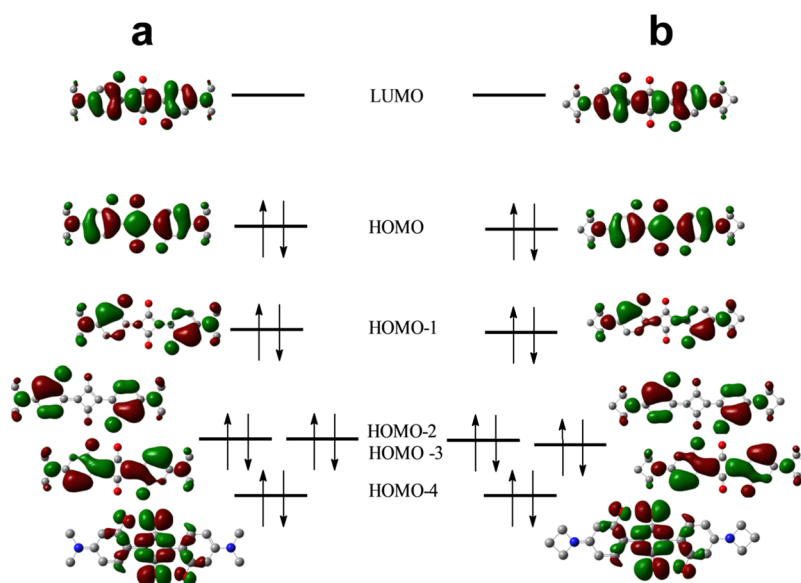


Figure 10. Main molecular orbitals of **1** (a) and **2** (b).

2. It is worth mentioning that the relatively small extinction coefficient $\epsilon^{\max} \approx 1.35 \times 10^5 \text{ M}^{-1} \text{ cm}^{-1}$ for new squaraine **1** is not explained by the employed theoretical approach. As follows from the data in Table 2, calculated transitions into the higher excited electronic states $S_0 \rightarrow S_n$ ($n = 2-5$) exhibit weak intensity ($f < 0.08$), and the different nature depends on the participating molecular orbitals. This prediction is in good agreement with the very weak absorption in the short wavelength range of the experimental IPA spectra of **1** and **2** (Figure 2, curves 1). The main frontier molecular orbitals of **1** and **2** (Figure 10) participating in the first 5 electronic transitions belong to HOMO type and only one to the LUMO type that is typical of symmetrical squaraines in the red-NIR spectral range.^{17,18} As follows from the data in Table 2 and Figure 10, there is only one $n \rightarrow \pi^*$ electron transition ($S_0 \rightarrow S_2$), which involves a lone electron pair at the corresponding oxygen atoms and exhibits negligible transition dipole moment. In general, the nature of the main electronic transitions of **1**

and **2** are similar to each other and mostly correspond to experimental results along with other calculated parameters.

4. CONCLUSIONS

A comprehensive investigation of the linear photophysical, 2PA, supperluminescence, and photochemical properties of new symmetrical squaraine derivative **1** was performed in liquid solutions at room temperature along with the quantum chemical analysis. Steady-state absorption, fluorescence, and fundamental excitation anisotropy spectra of **1** revealed relatively narrow absorption and emission bands (FWHM ≈ 30 nm), small Stokes shifts (<15 nm), high quantum yields ($\Phi_f \approx 0.8-1.0$), only one electronic transition $S_0 \rightarrow S_1$ in the main long-wavelength absorption band, and an angle of $\approx 25^\circ$ between the absorption, μ_{01} , and emission, μ_{10} , transition dipoles. All these parameters are similar to known reference squaraine **2** (and typical for other symmetric squaraines in the red-NIR spectral range^{18,27}), except for the value of the maximum extinction coefficient $\epsilon^{\max} \approx 1.35 \times 10^5 \text{ M}^{-1} \text{ cm}^{-1}$.

This value is noticeably smaller than that expected from the linear steady-state and time-resolved spectral data and theoretical prediction, and the origin remains unclear. The degenerate 2PA spectrum of **1** was obtained over a broad spectral range and a maximum cross-section of $\sigma_{2PA}^{\max} \approx 400$ GM was shown at ~ 800 nm, where the double-resonance excitation conditions can be expected.

The phenomenon of superluminescence emission was observed for **1** and **2** in polar DCM at room temperature under 1 kHz femtosecond pumping with a corresponding threshold energy ~ 0.12 to 0.15 μ J and is potentially important for the development of spectrally bright fluorescence labels for bioimaging. The results of DFT/TD-DFT based quantum-chemical calculations were in good agreement with the main spectroscopic parameters of **1** and revealed the nature of the main 1PA band. Linear photophysical properties, 2PA efficiency, superluminescence phenomenon in the red-NIR spectral range, and high photostability of azetidyl squaraine **1** suggest that it has promising potential for two-photon fluorescence bioimaging applications.

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Notes

The authors declare no competing financial interest.

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