

Plasmon enhanced nonlinear absorption and refraction

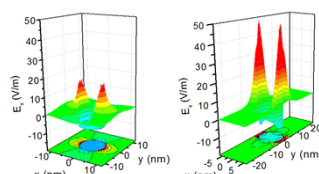
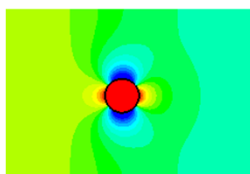
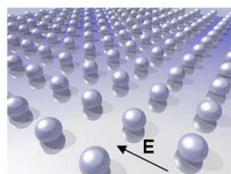
Pieter G. Kik

CREOL, The College of Optics and Photonics

University of Central Florida, Orlando, FL

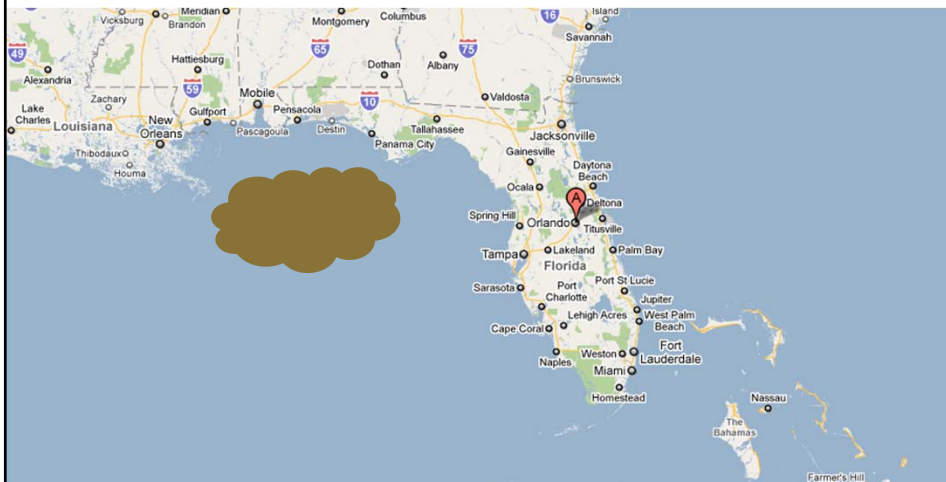
Calculations and simulations [Dana Kohlgraf-Owens](#)

Key idea: Use local fields to generate enhanced nonlinear optical response




ARO MURI *Engineered Multifunctional Nanophotonic Materials for Ultrafast Optical Switching*

Orlando, FL



UCF –outgrowing Google Maps

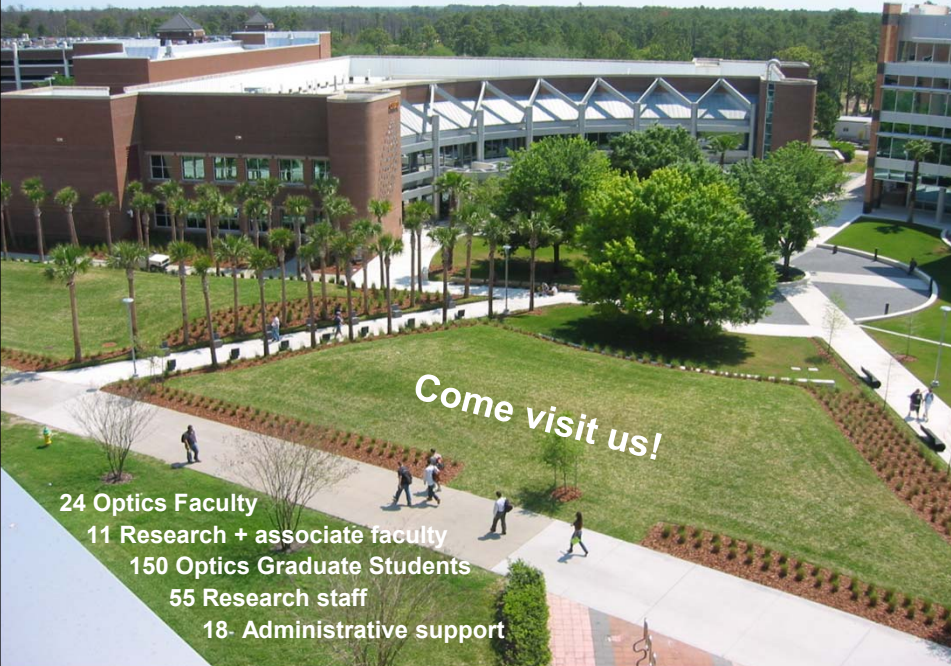


Founded in 1968
> 50,000 students
2nd largest in USA
> 1200 faculty
> 180,000 degrees awarded

**CREOL:
always evolving**

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Come visit us!

- 24 Optics Faculty
- 11 Research + associate faculty
- 150 Optics Graduate Students
- 55 Research staff
- 18- Administrative support

NPNO group – plasmon related work

Dielectric control of SP resonances

Strong coupling in surface coupled nanoparticle arrays (poster: Tuesday)

Near-field focusing using metal lenses

LSP enhanced nonlinear absorption

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Outline

1. Introduction: nonlinear absorption
2. Intermezzo (recovery from equations)
3. Plasmon enhanced NL absorption in dilute solutions
 - ⇒ surprise: complex enhancement
4. Figure of merit for NL absorption: role of plasmon induced damping
5. Optimization of NL absorption by nanostructure design
6. Summary

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Nonlinear refraction

Nonlinear refractive and absorptive materials are of interest for

- Nonlinear transmission (power management)
- Self-focusing
- All-optical switching
- Sensor protection

Holy grail: achieve low threshold for NL absorption in thin-film coating

Challenge: nonlinear optical response generally weak

⇒ Extremely large irradiance needed to achieve any significant effects

Question: can plasmon resonance on spherical particles be used to increase the nonlinear absorption performance in a thin sample

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Complex nonlinear refractive index

In the absence of 2nd order effects, polarization given by $P = \epsilon_0 \chi^{(1)} E + \chi^{(3)} E^3 + \dots$

Finite $\chi^{(3)}$ leads to (complex) change in n: $n = n_0 + \eta_2 I = n_0 + (\eta_2' + i \eta_2'') I$

Low irradiance


High irradiance

NL absorption

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Brief recap – propagation equations

In the absence of linear absorption, we have $\frac{dI}{dz} = -\beta I^2$



 nonlinear absorption coefficient $\beta = 2 k_0 n_2''$

This leads to nonlinear signal propagation: $I(z) = \frac{I_0}{1 + \beta I_0 z}$ (max $\frac{1}{\beta z}$)

Slab of NL absorptive medium \Rightarrow transmitted signal is limited

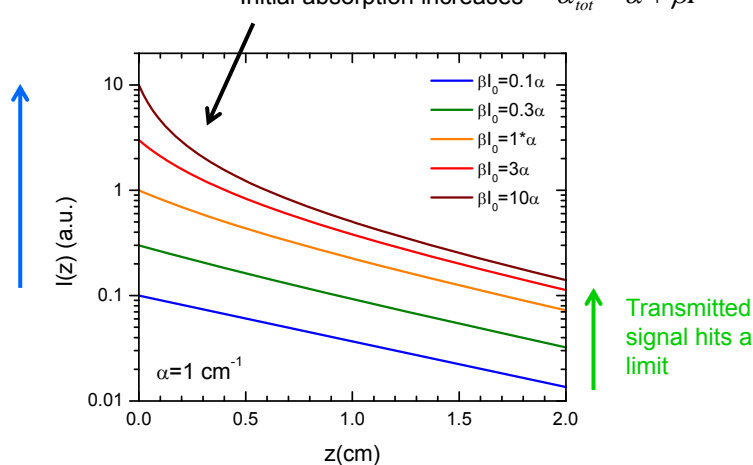
In the presence of linear absorption :

$$\frac{dI}{dz} = -(\alpha + \beta I)I \quad I(z) = \frac{1}{e^{\alpha_c L} - \frac{\beta_c I_0}{\alpha_c} (1 - e^{\alpha_c z})} I_0$$

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Example: nonlinear transmission

Initial absorption increases $\alpha_{tot} = \alpha + \beta I$



Challenge: irradiance required to achieve significant nonlinear absorption is high

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Effect of focusing on nonlinear absorption

Traditional solution: focusing incident radiation to increase irradiation

Focused lens lens

$I_{out} (W/m^2)$

$I_{in} (W/m^2)$

$\beta = 20 \text{ cm/GW}$
 $t = 1 \text{ cm}$

Threshold irradiance down
Max transmitted signal down

Practical issue: resulting system is large, relatively heavy

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Enhancing local fields using surface plasmon resonance

Known: large field strength \Rightarrow large nonlinear response

Clearly (..) : plasmon resonance and field enhancement will improve performance

ARO MURI "Engineered Multifunctional Nanophotonic Materials for Ultrafast Optical Switching" (CREOL / GATech / UArizona / Purdue)

Abbreviated proposal text: "Surround metal particle with nonlinear absorber \Rightarrow achieve enhanced NL absorption in small volume?"

note!

$$\frac{E_{surf}}{E_{inc}} = \frac{3\epsilon_m}{\epsilon_m + 2\epsilon_{host}}$$

20nm diameter Ag particle in air

First: evaluate using linear and 'nonlinear Maxwell Garnett theory'

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Intermezzo – finding Maxwell Garnett



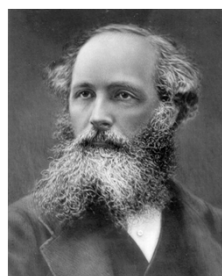
Who is **J. C. Maxwell Garnett** ?! Collaboration?

Google image search for a picture leads to many pictures of ...

James Clerk Maxwell (1831 - 1879) Famous for .. Maxwell's equations

BUT.. Effective medium paper on colored glasses published in **1904**..

Apparently we have: **James Clerk Maxwell** \neq **James Clerk Maxwell Garnett**



James Clerk Maxwell

Some internet searching reveals:

William Garnett, who was Maxwell's scientific demonstrator, was so impressed with Maxwell that he named his son James Clerk Maxwell Garnett



Mark I. Stockman Kik?

Outline



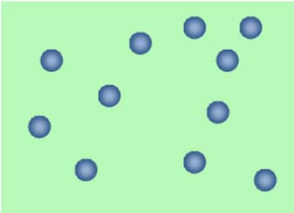
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Analytical calculations – Maxwell Garnet theory

Consider dilute solution of silver nanoparticles in CS₂

CS₂ (fs) $n_2=6 \times 10^{-15} \text{ cm}^2/\text{W}$ and $\beta=0.6 \text{ cm/GW}$
Data for 440nm – assumed constant

Assume linear Ag particle in NL host
 with fill fraction $f = 0.1\%$ and $d=10 \text{ nm}$



Nonlinear susceptibility of the composite is enhanced relative to host:

$$\chi_{\text{eff}}^{(3)} = g^{(3)} \chi_{\text{host}}^{(3)} \quad \text{and for NL index} \quad n_{2,\text{eff}} = g_2 n_{2,\text{host}}$$

Analytic expression for the enhancement factor:

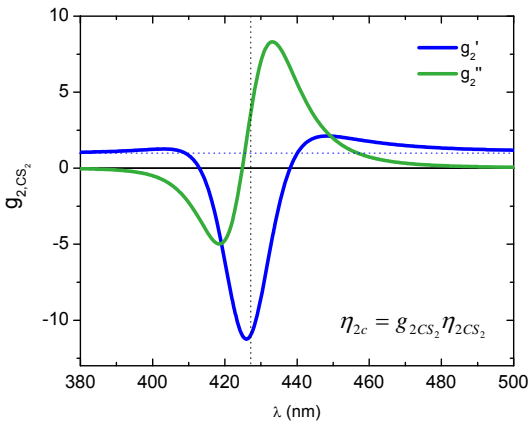
$$g^{(3)} \equiv 1 + \frac{f}{5} \left[8|\gamma|^2 \gamma^2 + 6|\gamma|^2 \gamma + 2\gamma^3 + 18(|\gamma|^2 + \gamma^2) \right]$$

M.I. Stockman et al.,
Phys. Rev. B 60, 17071 (1999)

Here γ is a complex field enhancement factor $\gamma = \frac{\epsilon_i - \epsilon_h}{\epsilon_i + 2\epsilon_h}$ ← plasmon resonance

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Nonlinear index enhancement: silver nanoparticles in CS₂

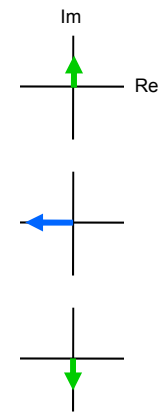


$\eta_{2c} = g_{2\text{CS}_2} \eta_{2\text{CS}_2}$

$\lambda \approx 440\text{nm}$

$\lambda \approx 425\text{nm}$

$\lambda \approx 415\text{nm}$



Good news: small amount of metal : $f=0.001$, but large enhancement factors

Surprise: at LSP resonance ($\sim 430\text{nm}$) \Rightarrow CS₂ becomes saturable absorber !

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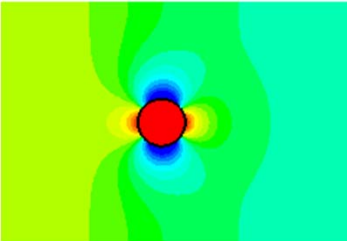
Nonlinear index enhancement: silver nanoparticles in CS₂

Origin of change in response:

$$P^{(3)} = \epsilon_0 \chi^{(3)} E^3$$

Locally enhanced fields in host:
90° phase delayed relative to incident light

$$P_{loc}^{(3)} = \epsilon_0 \chi^{(3)} E_{loc}^3$$

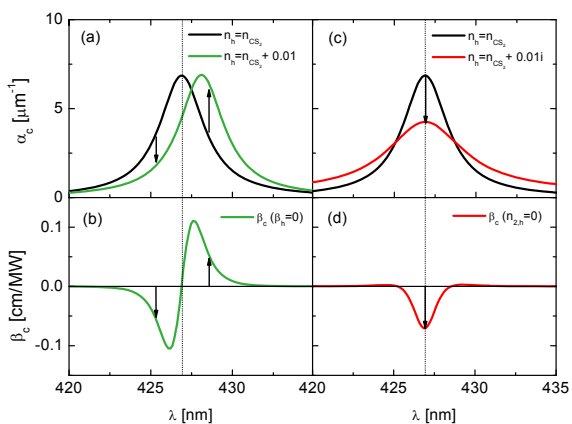
$$E_{loc} \propto iE_{in} \Rightarrow P_{loc}^{(3)} \propto -i\chi^{(3)} E_{in}^3$$


"Clearly plasmon resonance and field enhancement will improve performance .." - P.G. Kik, GRC 2010

Field enhancement does NOT simply enhance nonlinear absorption

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Interpretation: light-induced resonance changes

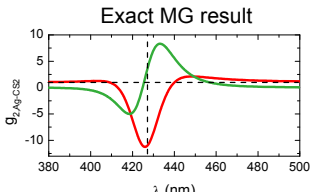


$$\sigma_{ext} = 9 \frac{\omega}{c} \epsilon_h' \frac{3}{2} V \frac{\epsilon_m''}{(\epsilon_m' + 2\epsilon_h')^2 + \epsilon_m''^2}$$

ε_H' up ⇒ resonance redshifts

ε_H'' up ⇒ resonance broadens

Exact MG result



Damping reduced 10x
large Δn_H and Δκ_H assumed

Wrong viewpoint: 'the enhanced local fields enhance the response of the host'

Correct: the nonlinear host polarizability affects the composite polarizability

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Trade-off: field enhancement vs. linear absorption

Second challenge: metal nanoparticles introduce linear absorption
 To evaluate quality of NL absorber, consider Figure of Merit $FOM = \beta / \alpha$

Field enhancement factor

$\gamma = \frac{\epsilon_{Ag} - \epsilon_{CS_2}}{\epsilon_{Ag} + 2\epsilon_{CS_2}}$

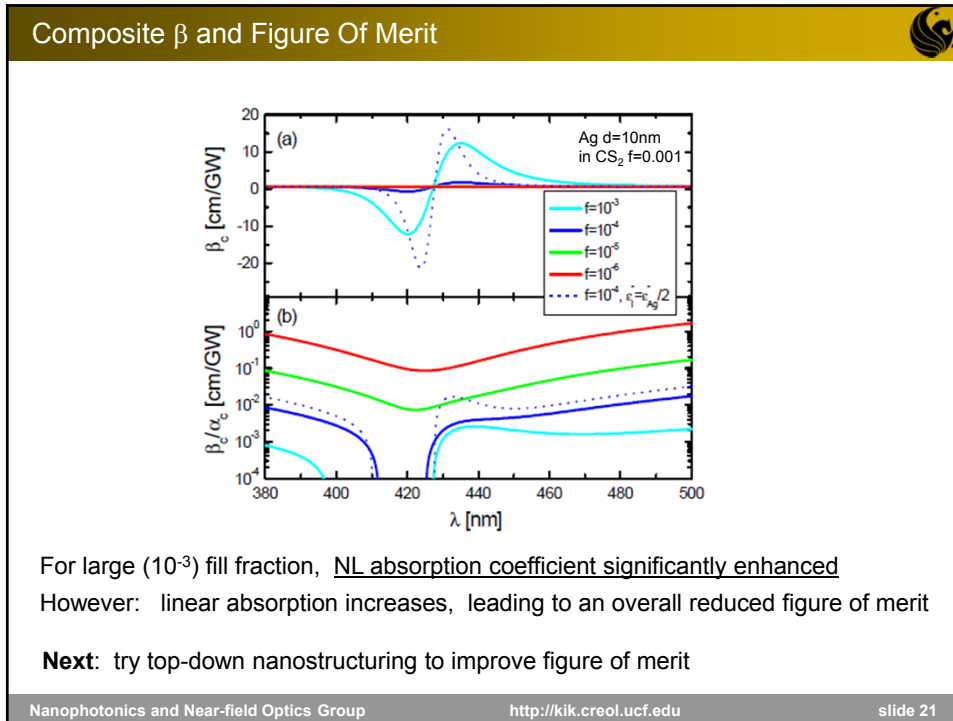
Ag d=10nm in CS₂, f=0.001
 resonance for $\epsilon_{Ag}' = -2\epsilon_{CS_2}'$

Linear absorption

$\epsilon_c = \left(\frac{1 + 2f\gamma}{1 - f\gamma} \right) \epsilon_h$

Max $|g^{(3)}|$ occurs at frequency for max linear absorption
Next: calculate figure of merit vs. fill fraction

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Structural engineering for optimum NLO absorption

Numerical simulations based on frequency domain simulations in CST Studio

Rectangular arrays of silver np (10nm diam, $f=3\%$) single layer (xy plane), $n_{\text{host}}=1.5$

Change inter-particle interaction by varying unit cell **aspect ratio**

Monitor geometry dependence of field enhancement and NLO response

Top view of geometries, incident E-field along x (horizontal), layer thickness 30nm

$L_x=12.5$ $L_x=15$ $L_x=24$ $L_x=38$ $L_x=46$

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Effective medium response through numerical integration

Dense arrays: no analytical formulas available

Based on numerically computed field distributions in different unit cells we calculate linear dielectric properties using

$$\epsilon_c(\omega) = \frac{\langle \epsilon(\omega, \vec{r}) \vec{E}(\omega, \vec{r})^2 \rangle_V}{\langle \vec{E}(\omega, \vec{r}) \rangle_V^2}$$

And nonlinear optical properties using

$$\chi_c^{(3)}(\omega) = \frac{\langle \chi^{(3)}(\omega, \vec{r}) |\vec{E}|^2 \vec{E}^2 \rangle_V}{|\langle \vec{E} \rangle_V|^2 \langle \vec{E} \rangle_V^2}$$

Note: χ values represents χ_{xxxx} , only valid along certain structural symmetry axes

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Definition of enhancement factor

Composite NL susceptibility can be written in terms of enhancement factors g:

$$\chi_c^{(3)} = f_{in} g_{in}^{(3)} \chi_{in}^{(3)} + f_h g_h^{(3)} \chi_h^{(3)}$$

Note: fill fraction no longer part of the enhancement factor. New factors given by :

$$g_j^{(3)} = \frac{\langle \vec{E}^2 | \vec{E}^2 \rangle_{V_j}}{\langle \vec{E} \rangle_V^2 \langle \vec{E} \rangle_V^2}$$

← V_j = volume of inclusion or of host (j=in or h)
← V = volume of unit cell

These represent *enhancement of $\chi^{(3)}$ contribution from host or inclusion, relative to the expected value based on a homogeneous E distribution*

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Simulation structure

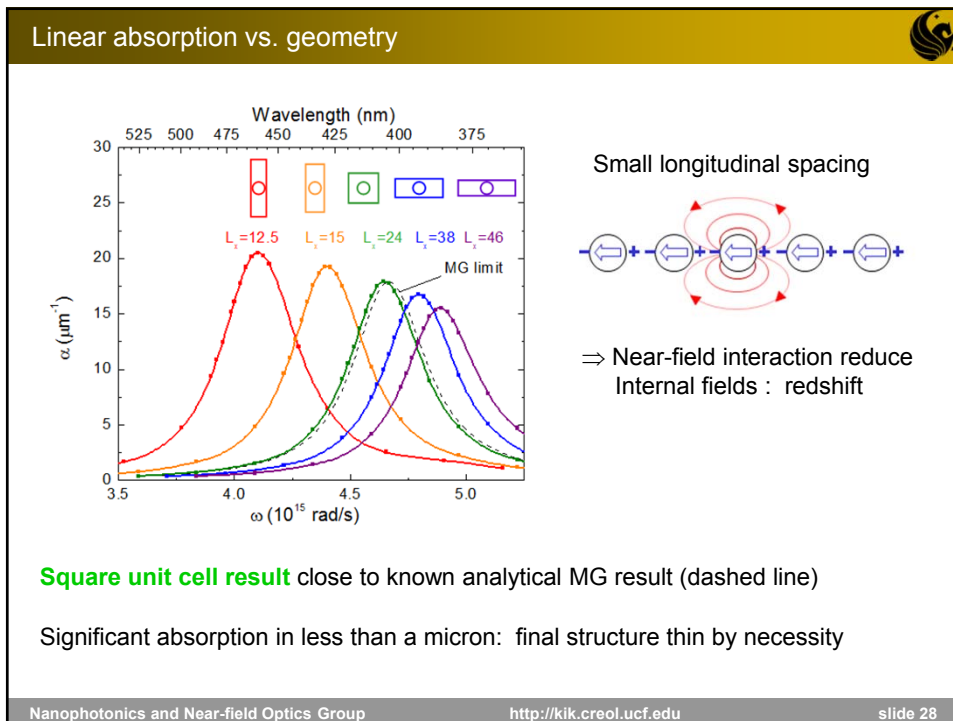
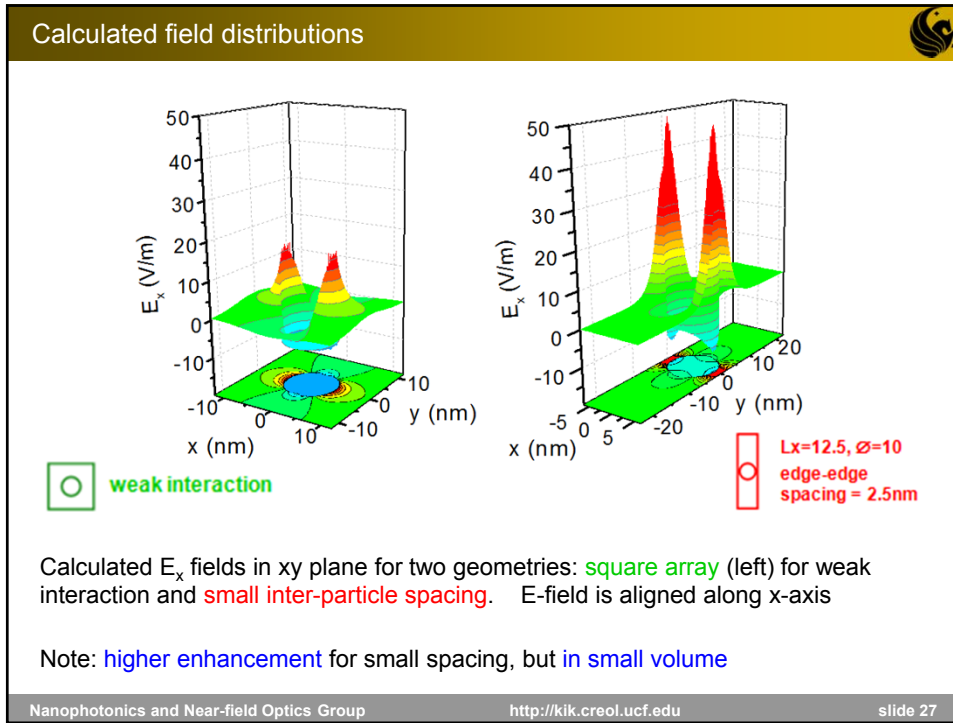
- Simulate 5 geometries ; **maintain fixed fill fraction $f=0.03$**
- 10 nm diameter Ag spheres in host index **$n=1.5$**
- Incident E field polarized along x direction

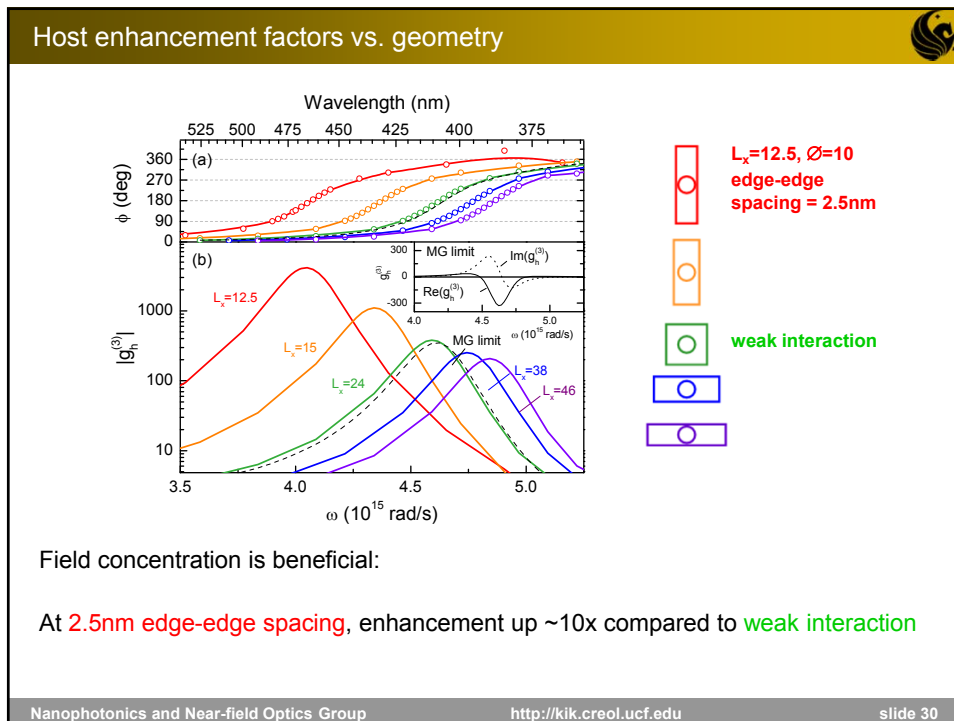
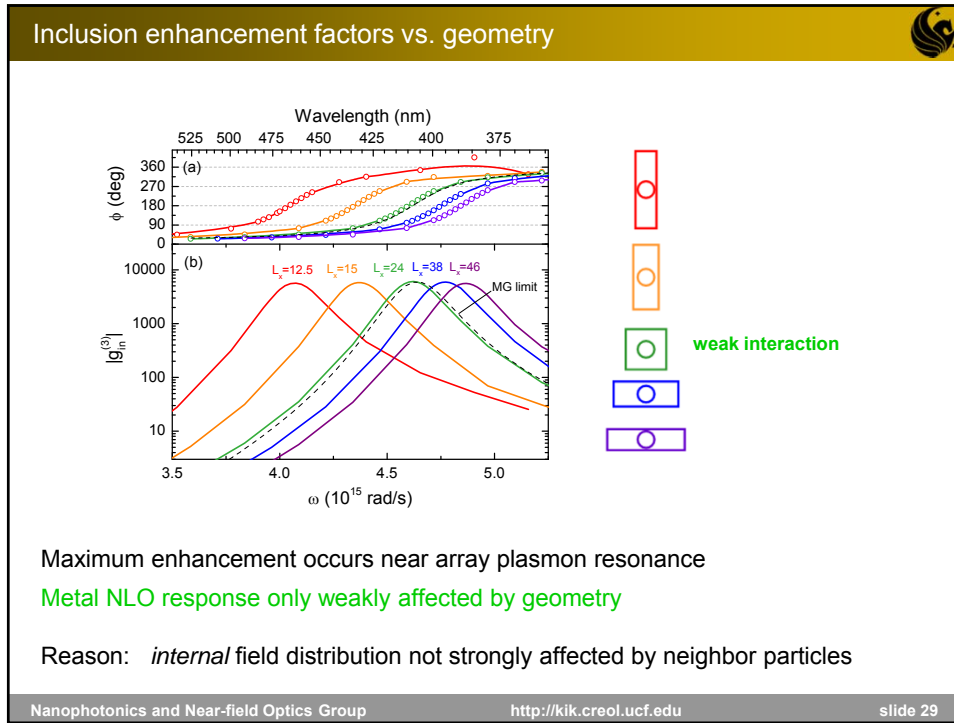
longitudinal interaction

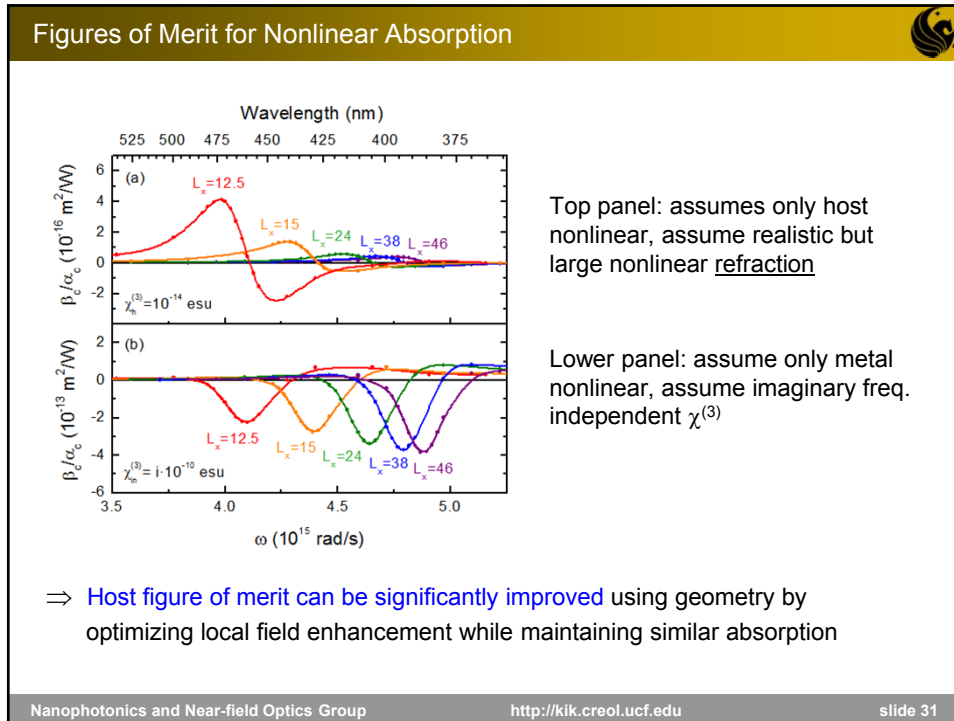
minimal interaction

transverse interaction

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Summary

Holy grail: achieve low threshold for NL absorption in thin-film coating

Dilute metal nanoparticle solutions:

- Small metal fill fractions generate large changes (>10x) in $\chi^{(3)}$ and n_2
- Enhancement is complex: NL absorptive host not best choice?
- Understand in terms of linear absorption spectra
- Figure of merit is reduced in presence of metal due to linear absorption, BUT thin film response is improved

Metal nanoparticle arrays

- Additional field enhancement at small spacing improves FOM

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