## 11.45 QMD4

## Debris -- free, droplet laser plasma sources in the EUV and soft X-ray ranges

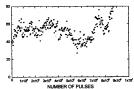
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Laser plasma short wavelength sources are finding increasing applications in science and advanced technologies. In the development of Extreme Ultraviolet Lithography (EUVL), the leading technology to replace optical lithography, the source of choice is a high repetition-rate (> 1kHz) laser-plasma, capable of providing at least 7 W or Tradiation within a 3% spectral bandwidth at 13 mu Laser plasma exhallow of providing at least 7 W or Tradiation within a 3% spectral bandwidth at 13 mu Laser plasmas make excellent pulsed x-ray sources in the so-called "varter window" (2.5 – 45 nm) for nanometer-resolution x-ray microscopy of biological organisms in their natural state. And plasmas produced by femiosecond lasers are now a source of high-energy x-ray radiation for x-ray diffraction and radiography. A common requirement of these applications, that there be no collateral particulate emission from the plasma, has lead to the development of the liquid droplet source, introduced in 1993, as a bright source in the soft x-rays and EUV radiation. Our previous work with a 100 kHz water droplet system generated 13 nm and 11 6 nm line emission from a Li-like Oxygen plasma produced by 10 ns duration, 10 Hz Nd: YaG laser plasts at -10<sup>12</sup> W/cm².

We now present a detailed quantitative study of this source, performed with a 100 Hz laser, that characterizes the radiation efficiency and the long-term operation. The results show that the droplet laser plasma source comes close to satisfying all the near-term needs of EUV. In particular we demonstrate an overall conversion efficiency of laser light to 13 nm emission within the required spectral bandwidth in excess of 06%, comparable to any other existing source at this wavelength. In addition we have performed an exhaustive examination of the long-term effects of plasma emissions on the reflectivity of multi-layer mirrors, exposed to this source. This included mirrors, exposed to this source. This included mirrors, exposed to this source. This included mirrors, exposed to this sourc



- L. Ryssell & H. Hertz, Opt. Comm. 103, 105, (1993): M. Richardson, K. Gabel, F. Jin & W. T. Silfvast, <u>Proc. OSA Top. Mtg. Soft x-ray Projection Lithogonalth</u>, OSA, Washington DC, Vol. 18, pp. 156-162, (1993)
  F. Jin & M. Richardson, Appl. Opt. 3-m, 5750, (1995): M. Richardson, D. Torres, C. DePriest, F. Jin & G. Shimkaveg, <u>Opt.</u> Comm., 145, 106 (1998)

## 12.00 QMD5

## Laser Operation with Two Controlled Transverse Modes

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In a laser resonator operating with many transverse modes, the emerging output beam quality is relatively poor. Improvement of the beam quality is typically obtained by inserting an aperture inside the resonator so as to reduce the effective radius of the gain medium, until only the optimal fundamental mode of Gaussian shape exists. Unfortunately, this results in a significant reduction of the output power. It is possible to operate a laser with a single high-order transverse mode by applying intra-cavity elements, such as apertures or phase elements [12]. Such high-order modes exploit a relatively large volume of the gain medium, so the output power is relatively high. However, the intensity distribution of a single high-order mode usually consists of some low-intensity regions, resulting in inefficient utilization of the gain medium, with respect to the multi-mode operation. Here we demonstrate a novel resonator configuration in which it is possible to select two different transverse modes, each of which has a specific polarization, and can be manipulated esperately. The intensity distributions of these modes can be chosen to be complementary, i.e. the peaks of the first mode would fall on the valleys of the second mode. Thus, the gain medium can be exploited more efficiently. This results in an increase in the total output power, yet with better beam quality than with the multi-mode operation.

A representative resonator configuration in for a laser operation with the fundamental TEMoo mode (selected with an aperture) and the Laguerre-Gaussian TEMoo mode (selected with an aperture) and the Laguerre-Gaussian TEMoo mode (selected with an aperture) and the Laguerre-Gaussian TEMoo mode (selected with an aperture) and the Laguerre-Gaussian TEMoo mode (selected with an aperture) and the Laguerre-Gaussian TEMoo mode (selected with an aperture) and the Laguerre-Gaussian TEMoo mode (selected with an aperture) and the Laguerr

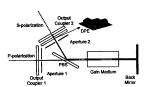




Fig 1: Resonator configuration

Fig. 2: Output beam intensity distribution

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 R. Oron, Y. Danziger, N. Davidson, A.A. Friesem, E. Hasman, Opt. Commun. 169, 115 (1999)